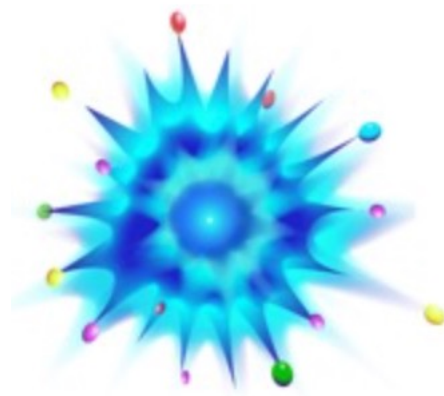




Probing the Nature's Primordial Fluid

Volodymyr Vovchenko (University of Houston)

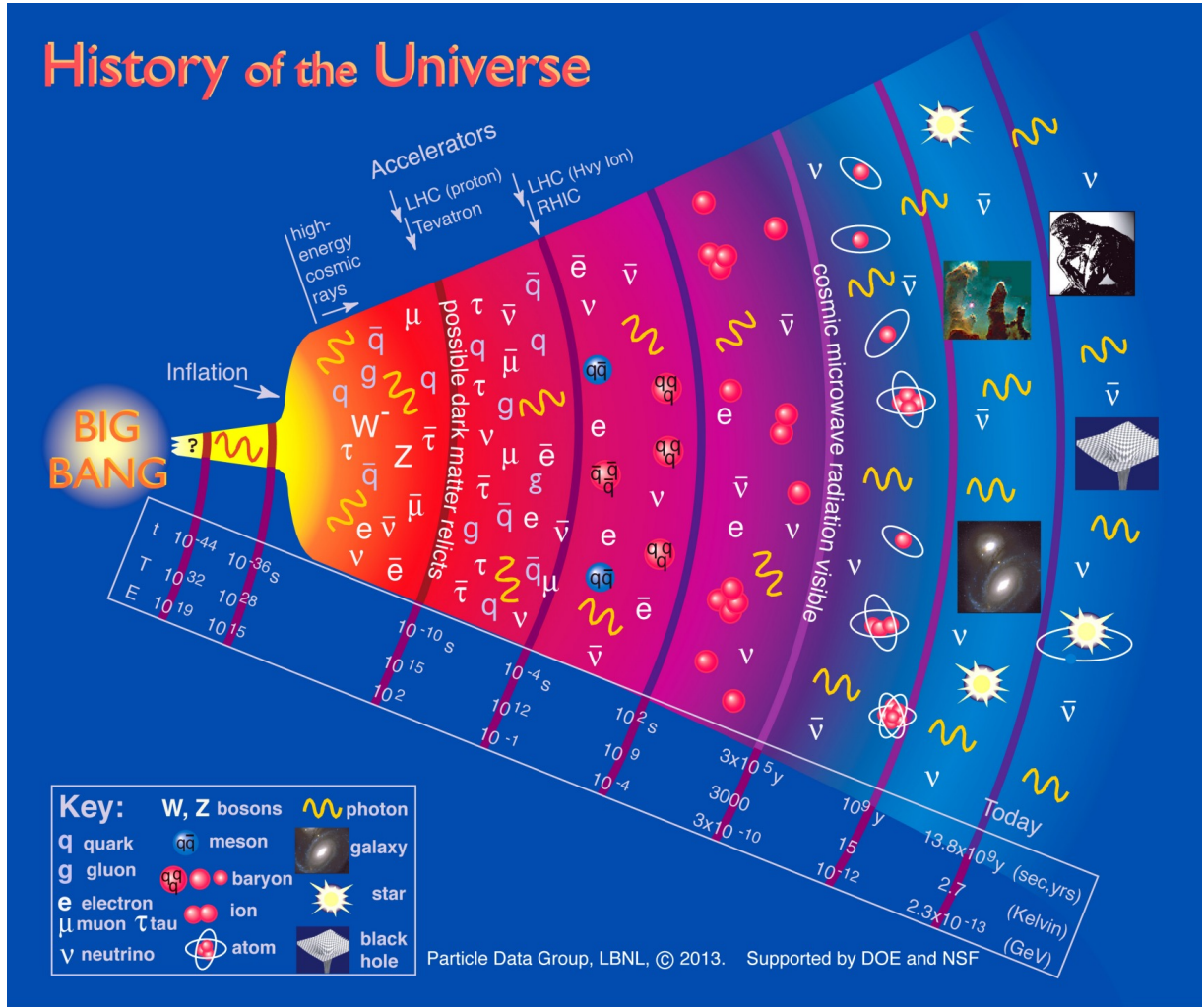
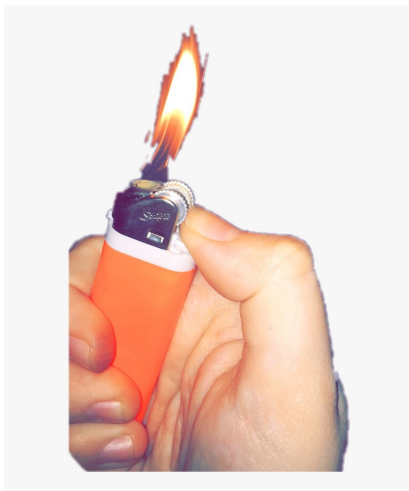
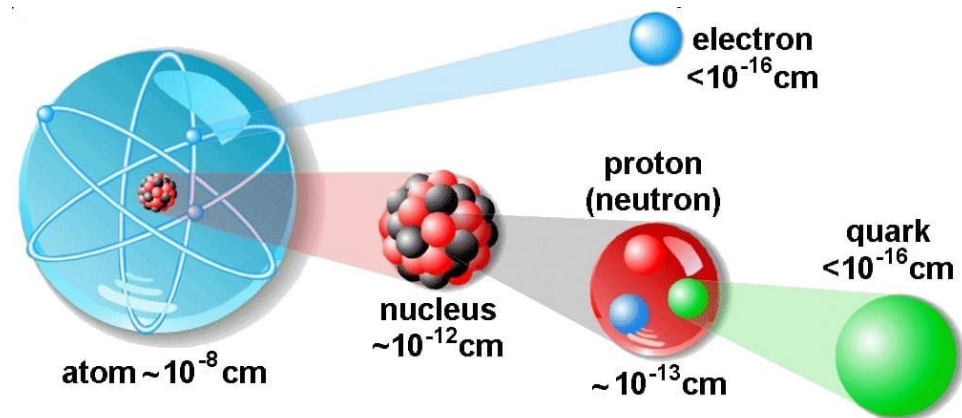
vvovchen@central.uh.edu – <https://vovchenko.net>

QuarkNet Summer Workshop 2025 – Rice University

June 10, 2025



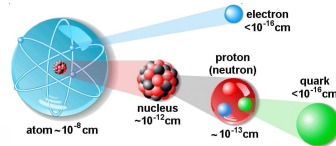
Structure of matter



Extreme states of matter



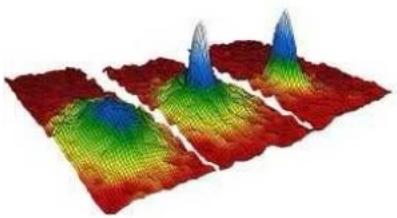
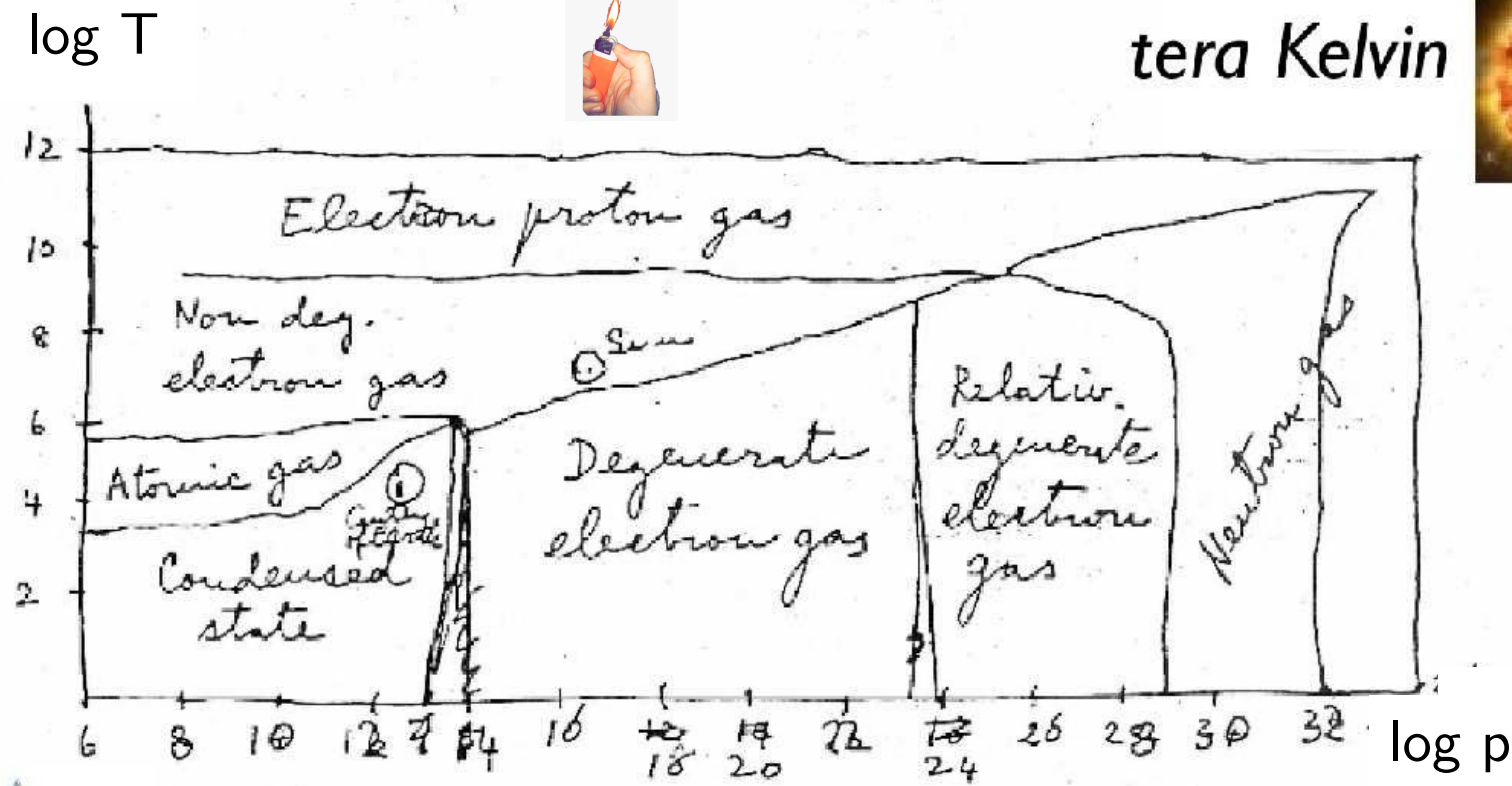
Fermi 1953



tera Kelvin



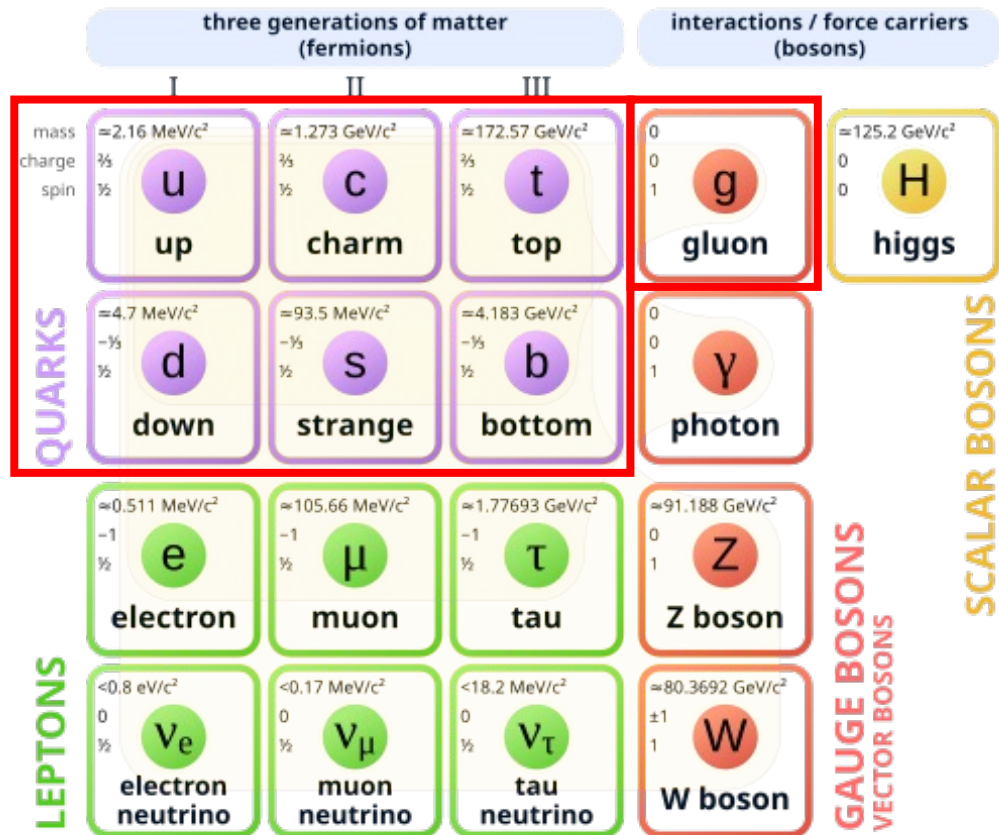
1 MeV $\sim 10^{10}$ K



Matter in unusual conditions

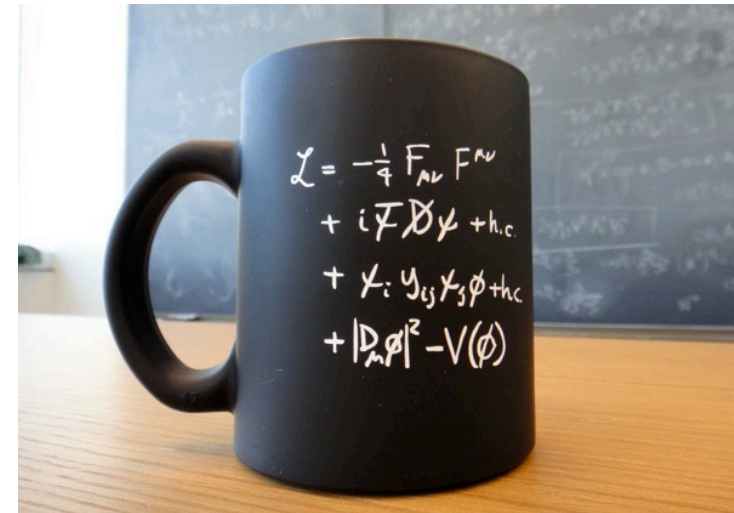
Standard model

Standard Model of Elementary Particles



mid-1970s

Predicted the top quark (disc. 1995), the tau neutrino (disc. 2000), and the Higgs boson (disc. 2012)



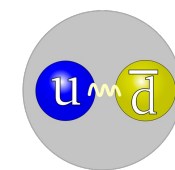
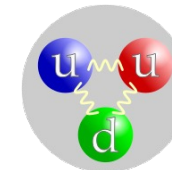
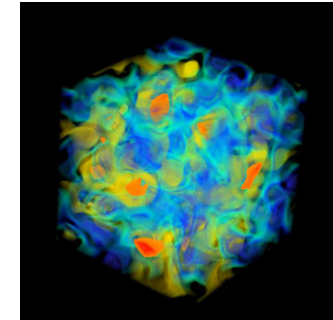
- Strong (nuclear) force
 - Quarks and gluons, short-range
- Electromagnetic force
 - Charged particles (Coulomb), long-range
- Weak force
 - Fermions (incl. neutrinos)

Strongly interacting matter

- Theory of strong interactions: *Quantum Chromodynamics* (QCD)

$$\mathcal{L} = \sum_{q=u,d,s,\dots} \bar{q} \left[i\gamma^\mu (\partial_\mu - igA_\mu^a \lambda_a) - m_q \right] q - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}$$

- Basic degrees of freedom: quarks and gluons that carry color charge
- At smaller energies confined into baryons (qqq) and mesons ($q\bar{q}$)



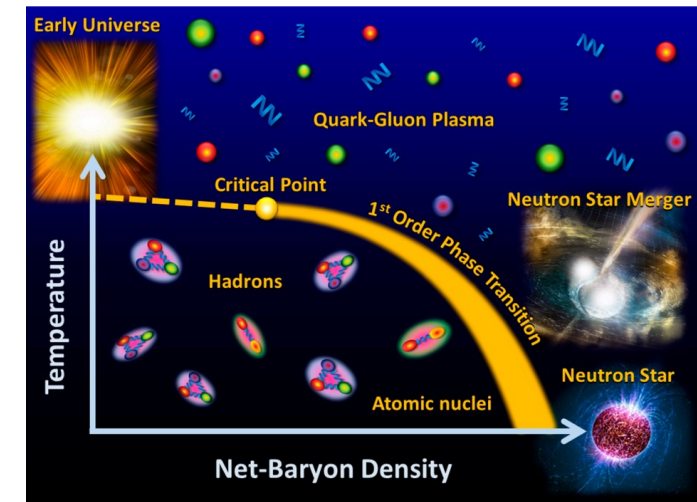
Scales

- Length:** 1 femtometer = 10^{-15} m
- Temperature:** $100 \text{ MeV}/k_B = 10^{12}$ K

Where is it relevant?

- Early Universe
- Astrophysics: Neutron star (mergers)

Studied in **laboratory** with heavy-ion collisions

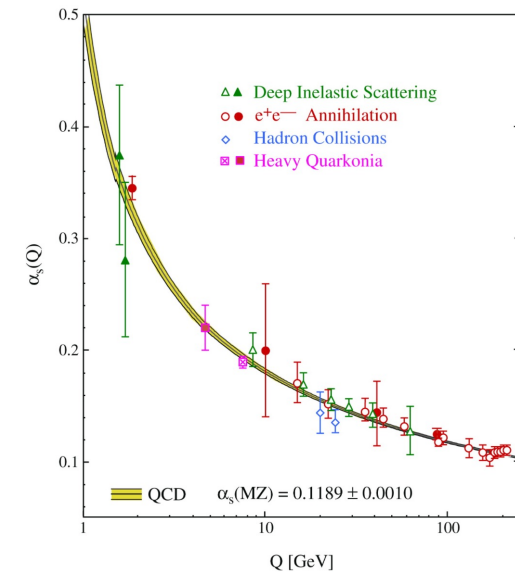


QCD features and emergent phenomena

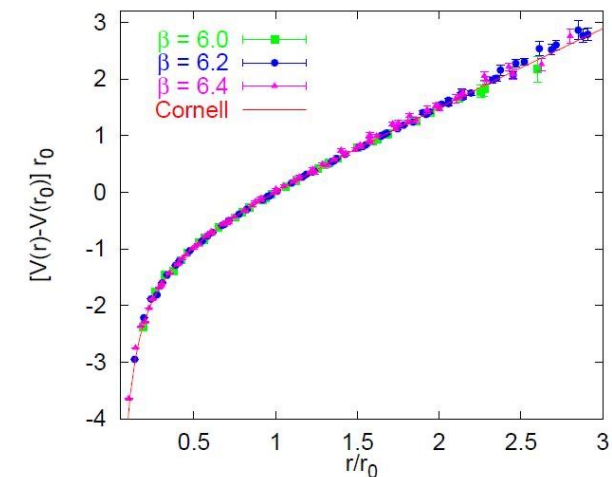
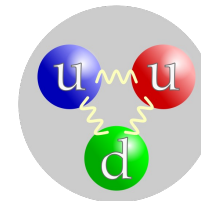
- Asymptotic freedom Gross, Politzer, Wilczek (1973)
 - Interaction becomes *weaker* at high energies/small distances
 - Theory is in perturbative regime at small distances
- Hadrons (confinement)
 - No free quarks or gluons ever observed
 - They must form composite, color-neutral objects – the hadrons
 - Proton (uud) and neutron (udd)
 - No small parameter makes the theory virtually untractable ☹
- Dynamical mass generation
 - Proton (uud) mass is $m_p = 938 \text{ MeV}/c^2$ but $m_u + m_u + m_d \sim 15 \text{ MeV}/c^2$
 - >95% of proton's mass from QCD, only <5% is from Higgs



2004



S. Bethke



G. Bali

- Forces usually get weaker with distance



Confinement

- Forces usually get weaker with distance

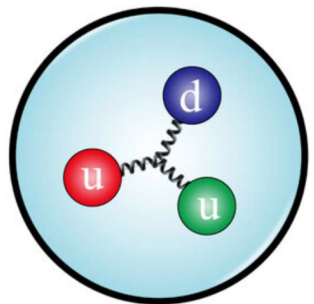

charged nucleus

can easily separate constituents

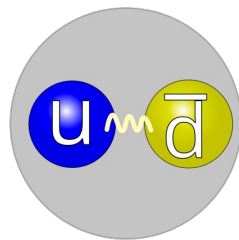
$$F \sim 1/r^2$$

 **electron**

- Quarks in QCD are confined



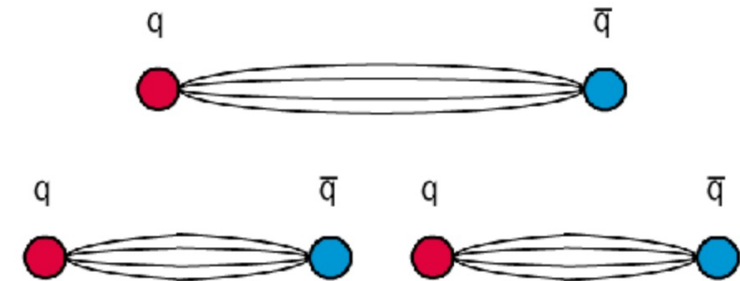
**“white” proton
(confined quarks)**



**“white” pion
(confined $q\bar{q}$)**

$$F \sim r$$

- Try to separate quarks

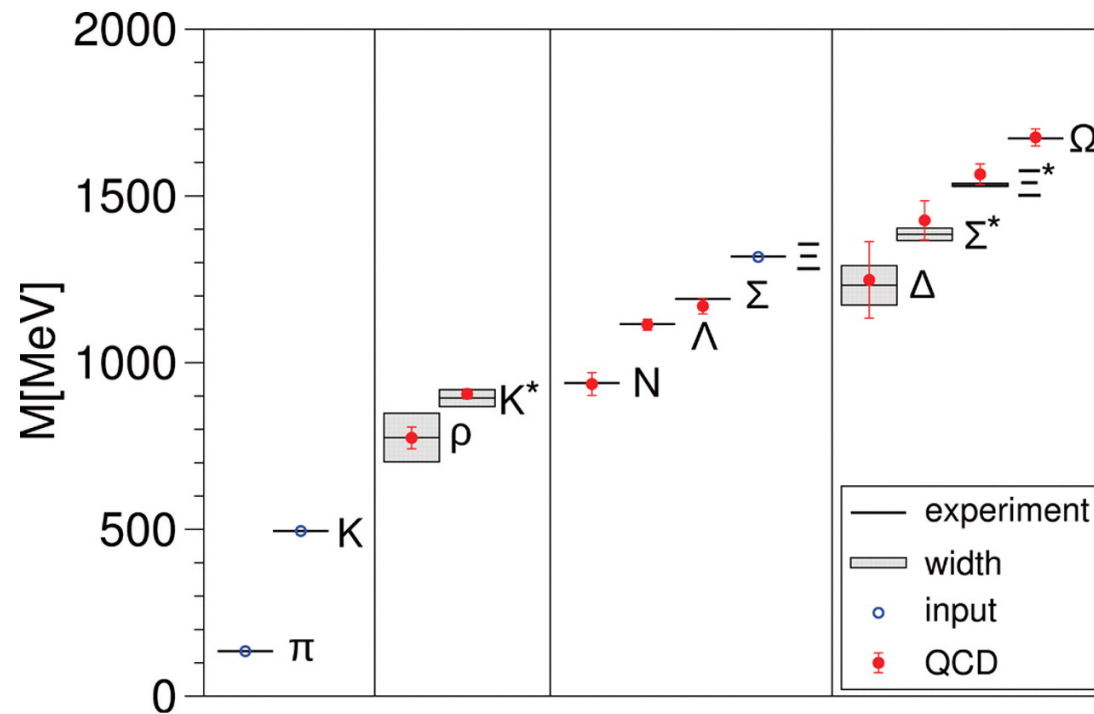


**quark-antiquark pair
created from vacuum**

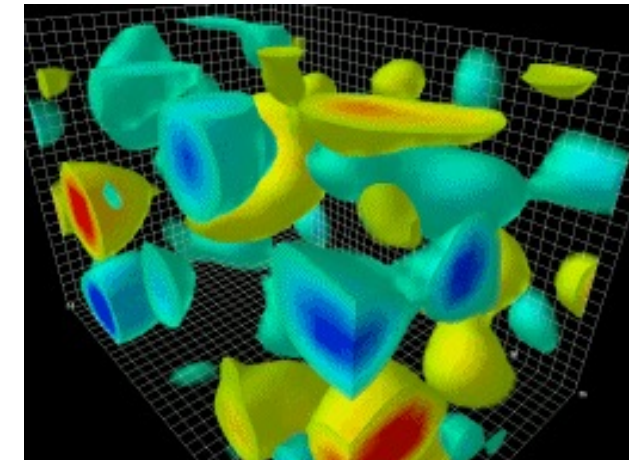
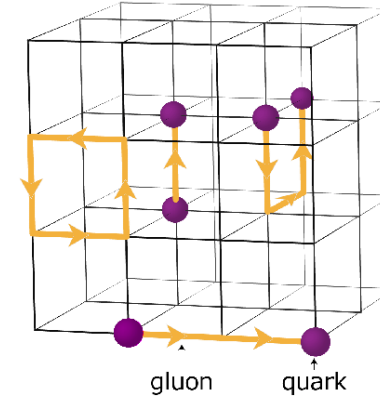
Non-perturbative methods

First-principle tool: Lattice QCD – a brute force Monte Carlo solution on a discretized space-time grid

Ab-initio calculation of hadron masses



BMW Collaboration, Science 322, 1224 (2008)

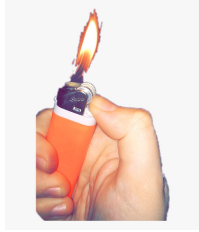


© CSSM, University of Adelaide

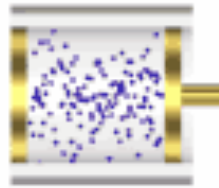
Remarkable agreement of QCD with the experiment

From confinement to deconfinement

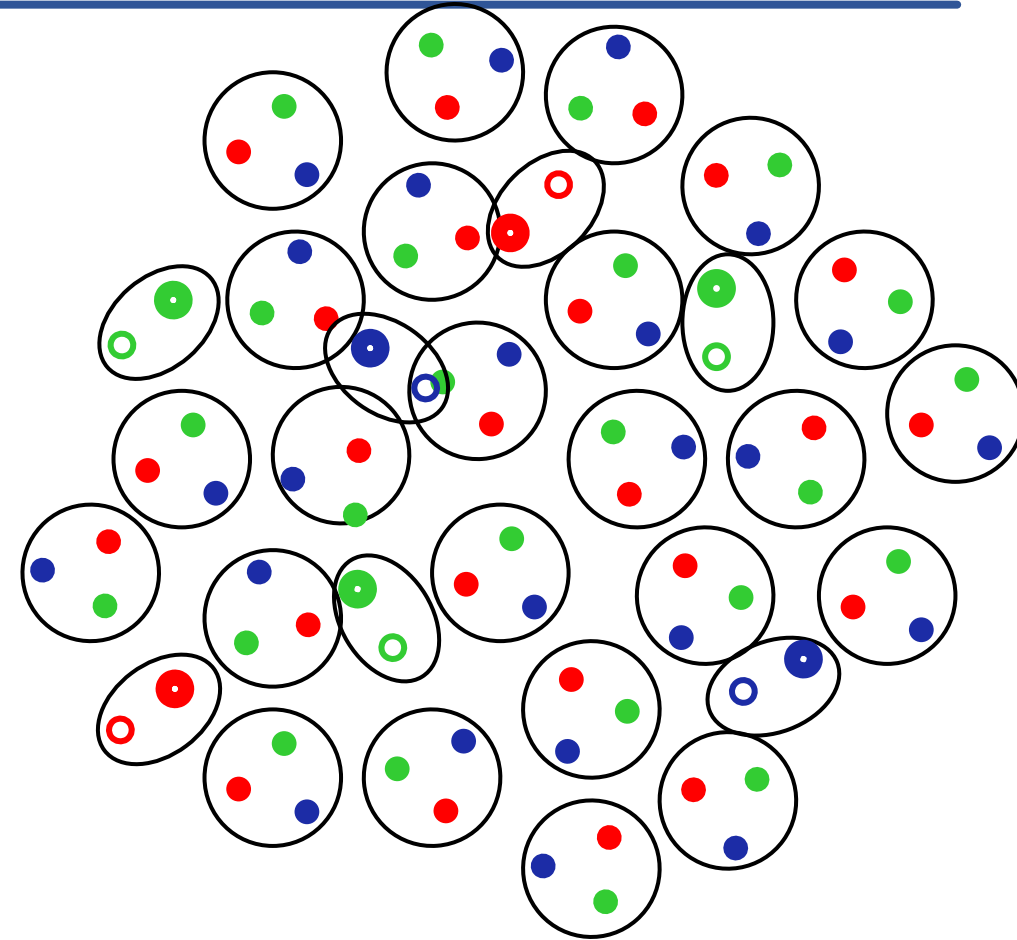
1. heating



2. compression



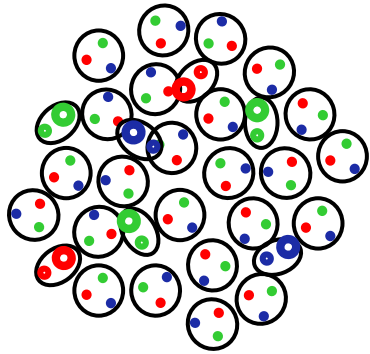
→ *deconfined color matter*
(*quark-gluon plasma*)



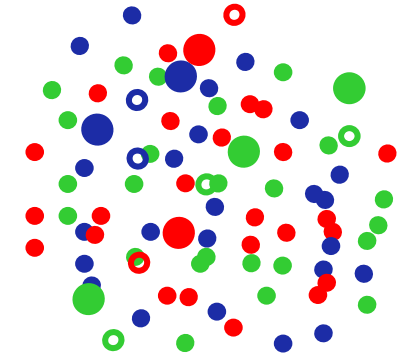
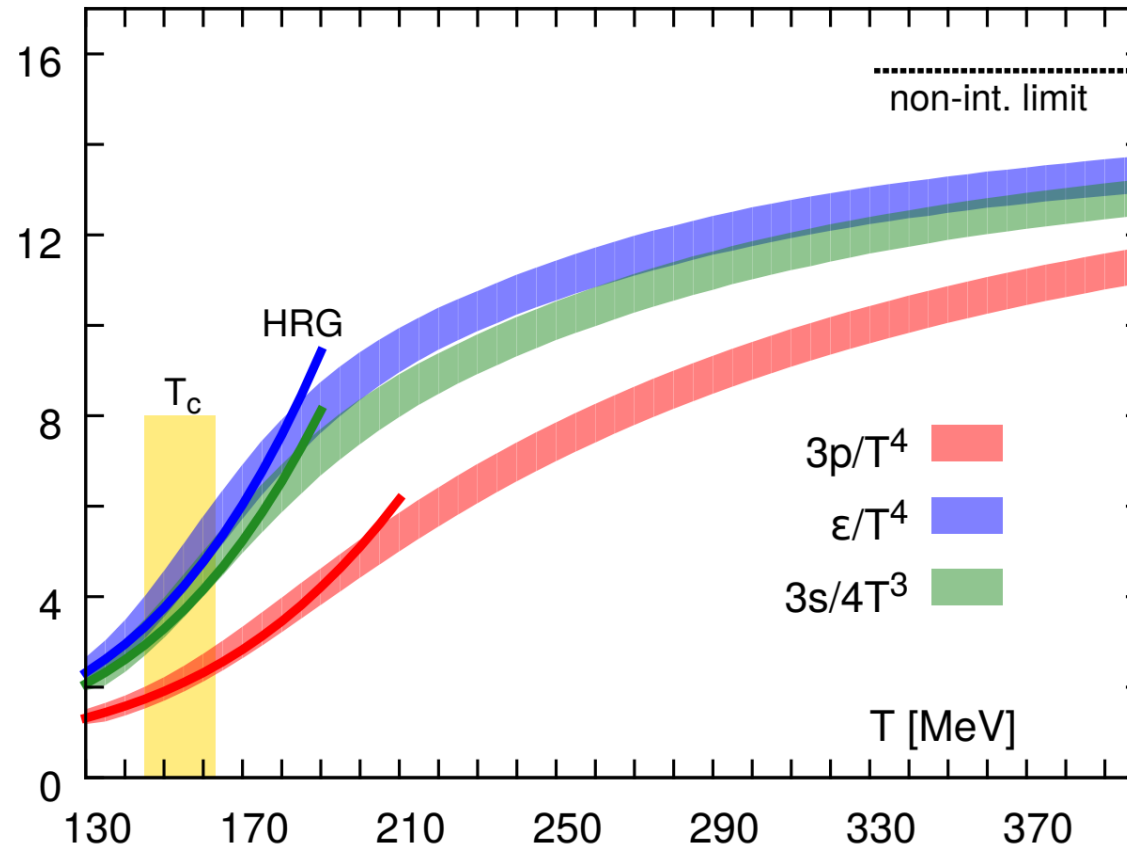
Quark-Gluon Plasma
(deconfined!)

Deconfinement transition from lattice QCD

$$\frac{\varepsilon}{T^4} \sim \# \text{ of d. o. f.}$$



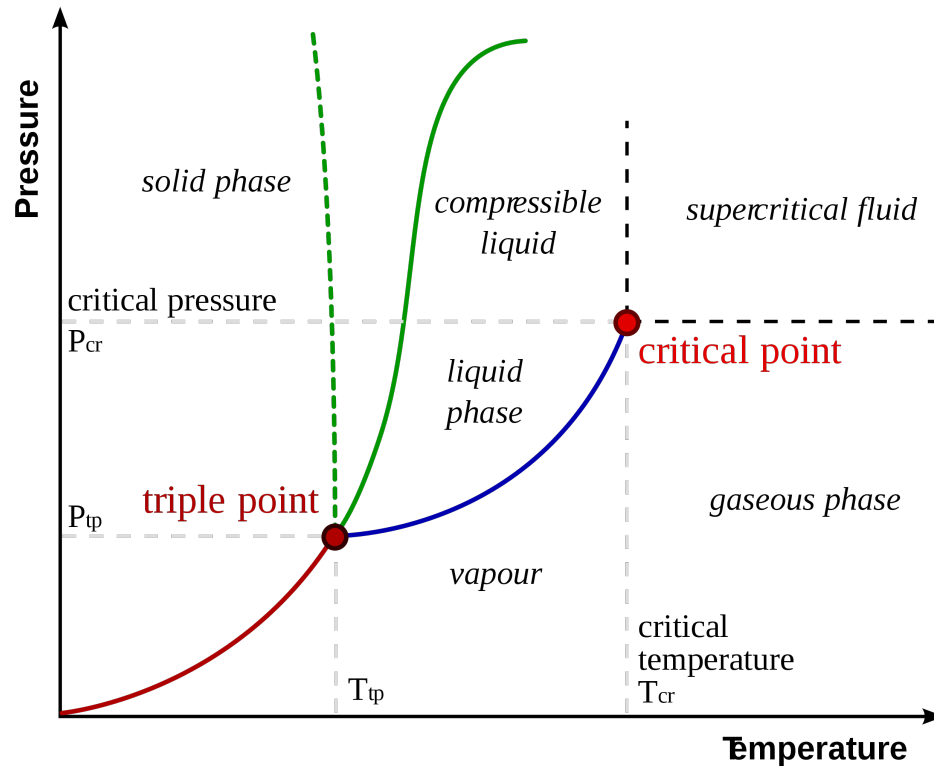
confined:
few d.o.f.



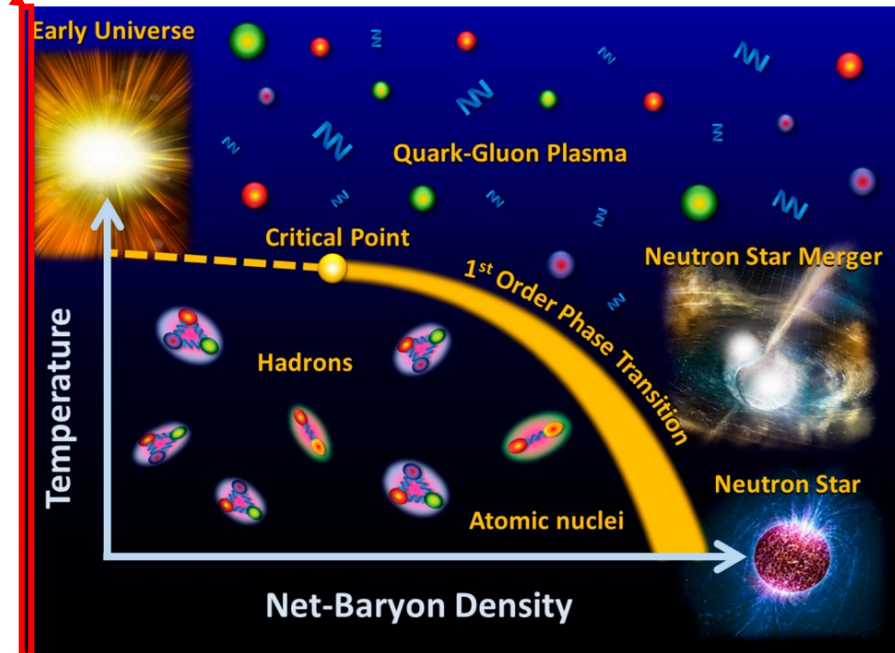
deconfined:
many d.o.f

Crossover transition at $T_c \approx 155 \text{ MeV} \approx 2 \cdot 10^{12} \text{ K} \approx 130,000 \cdot T$ [Sun's core]

QCD phase diagram



lattice QCD

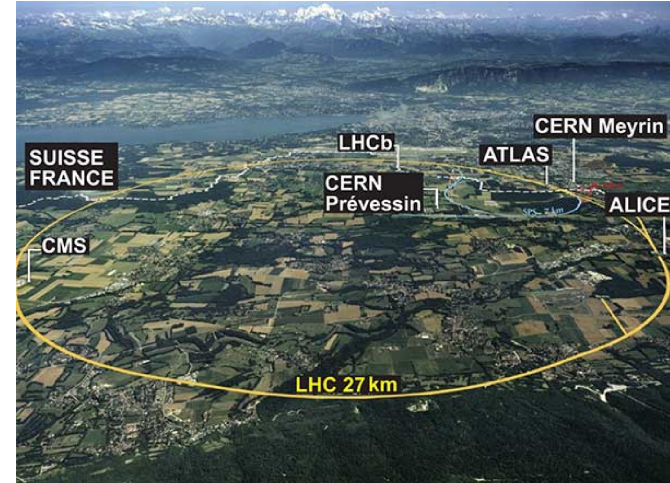


- Finite baryon densities inaccessible due to the sign problem
- Possible existence of a phase transition is a conjecture
- Laboratory: **heavy-ion collisions** – test of QCD and a tool unveil its many properties

QCD laboratories (~1980-...)



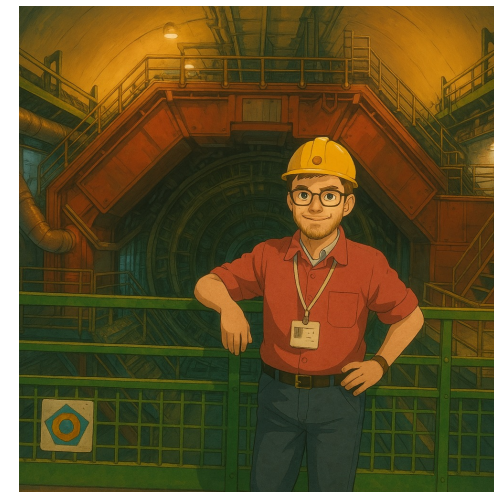
RHIC@BNL (2000-), 200 GeV, 99.995% speed of light



LHC@CERN (2010-), 5020 GeV, 99.9999991% speed of light

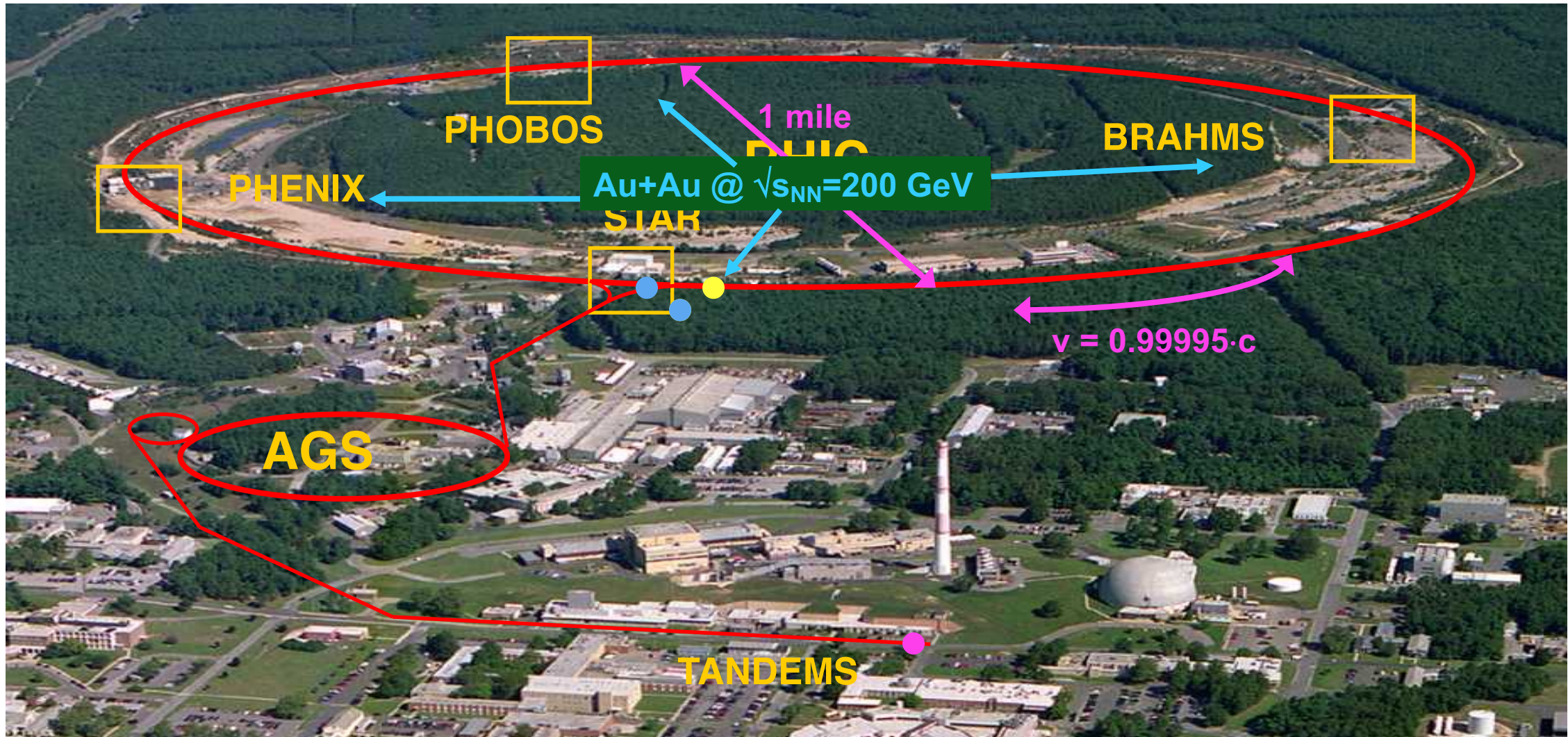


FAIR@GSI [2026(?)-], 5 GeV (less is more)



ALICE detector

Relativistic Heavy Ion Collider (RHIC)



Relativistic heavy-ion collisions

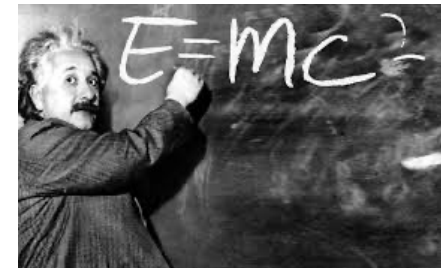
Pb + Pb, $E_{\text{kl}} = 158.0 \text{ A GeV}$

$b = 0.0 \text{ fm}$

Time: -19.50 fm/c



UrQMD model simulation



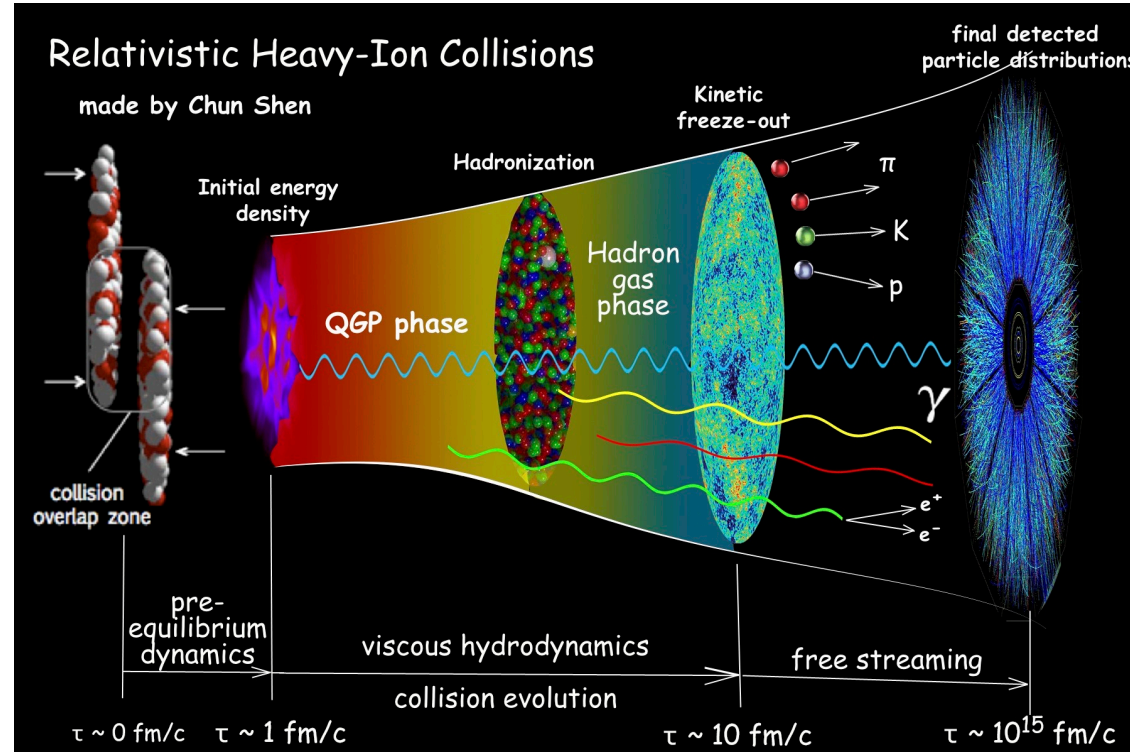
Relativistic collision energy
leads to particle production

Length: 10^{-15} m

Time: 10^{-22} s

Temperature: 10^{12} K

Relativistic heavy-ion collisions – “Little Bangs”



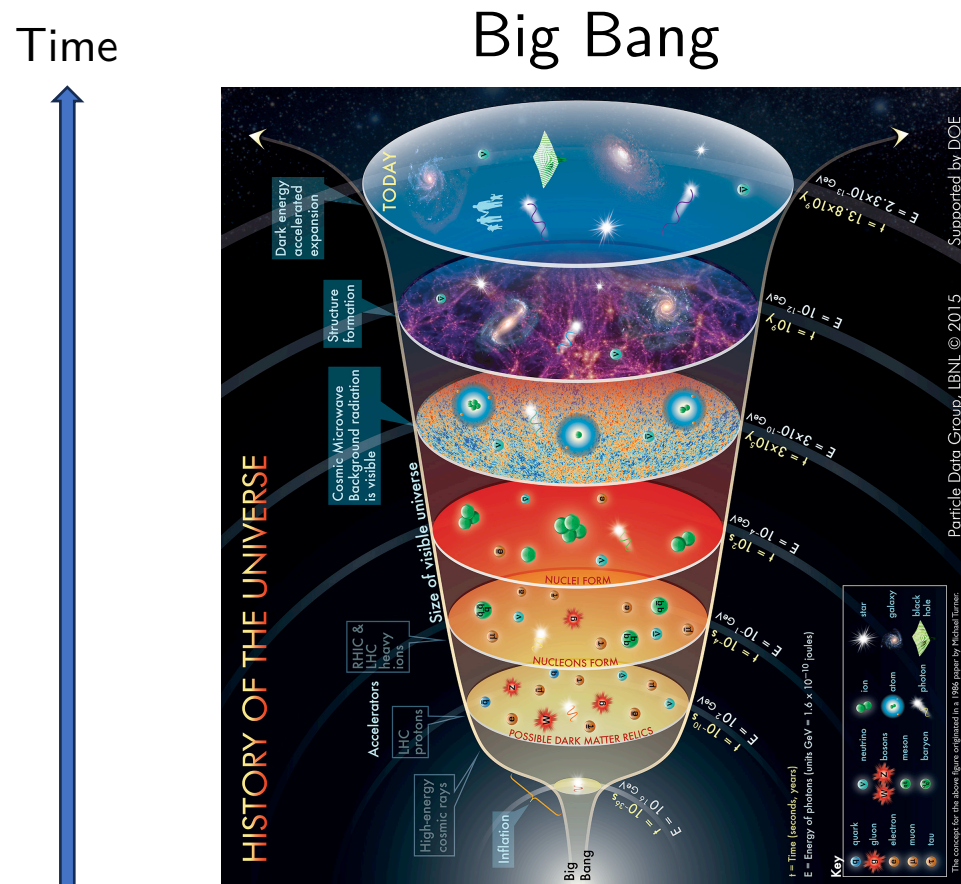
Control parameters

- Collision energy $\sqrt{s_{NN}} = 2.4 - 5020$ GeV
- Size of the collision region

Measurements

- Final hadron abundances and momentum distributions

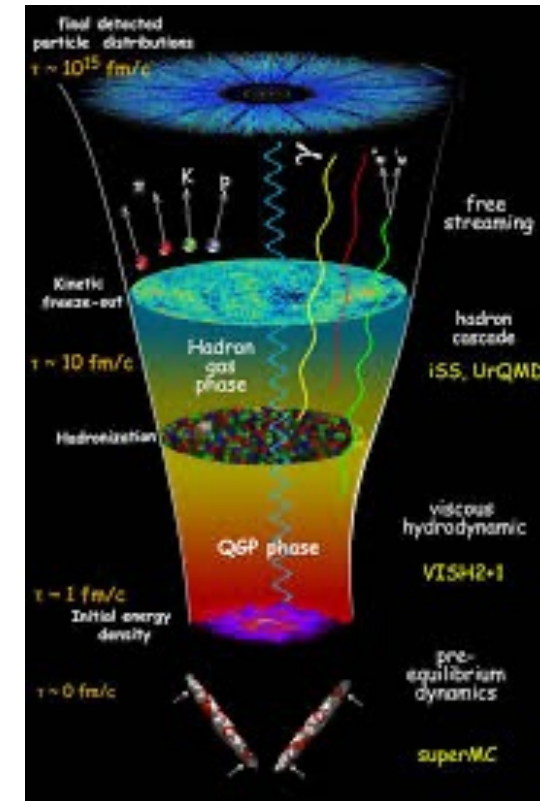
Big Bang vs Little Bang



1 event

Slow expansion, long-lived

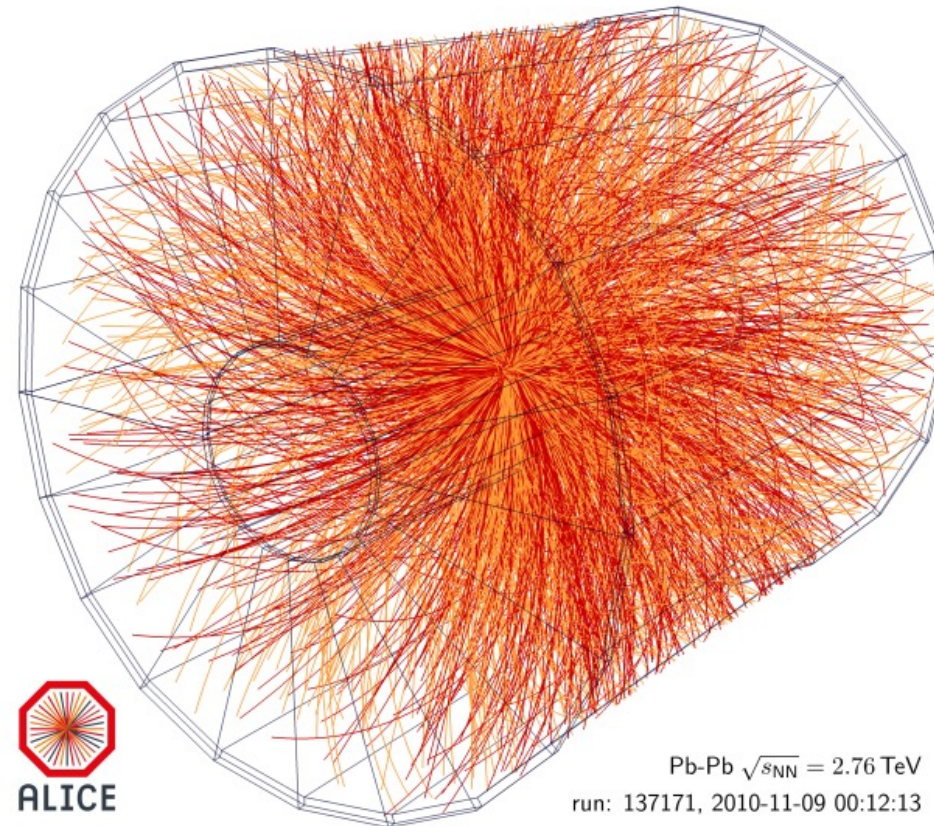
Heavy-Ion Collision



Millions of events

Rapid expansion, short-lived

Fireball in heavy-ion collisions



Event display of a Pb-Pb collision in ALICE at the LHC

Thousands of particles created in relativistic heavy-ion collisions



Apply concepts of statistical mechanics

Particle production in heavy-ion collisions

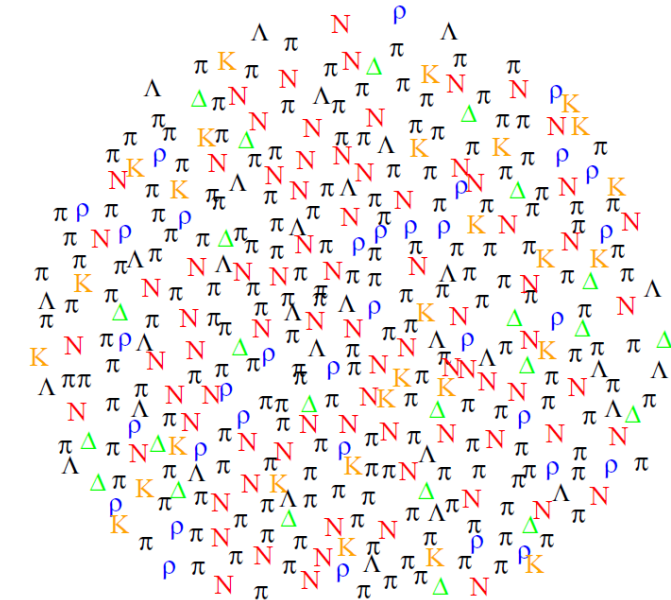
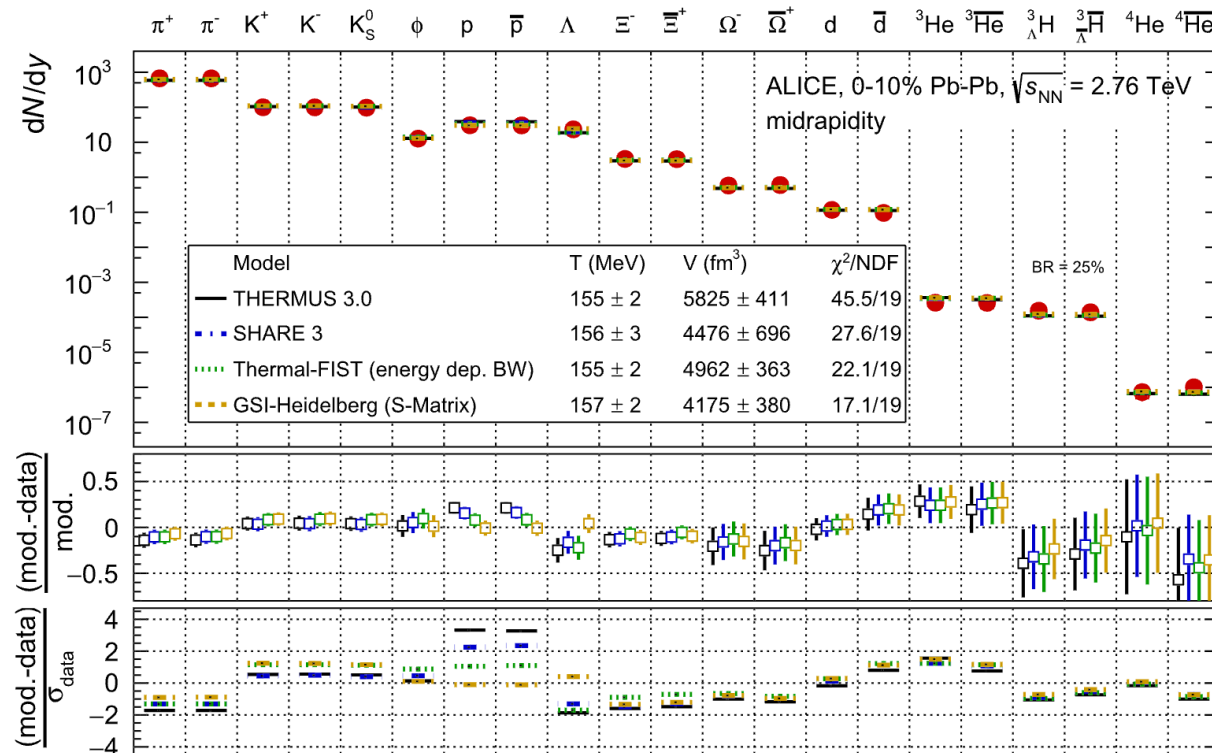
Ideal gas law (E. Clapeyron, 1834)

$$P_i V = N_i k_B T \quad (+ \text{feeddown})$$

is the simplest model of particle production

$$N_i = \frac{d_i V}{2\pi^2} \int dk k^2 \left[1 \pm \exp \left(\frac{\sqrt{k^2 + m_i^2} - \mu_i}{T} \right) \right]^{-1}$$

Bose-Einstein & Fermi-Dirac, 1924-1926

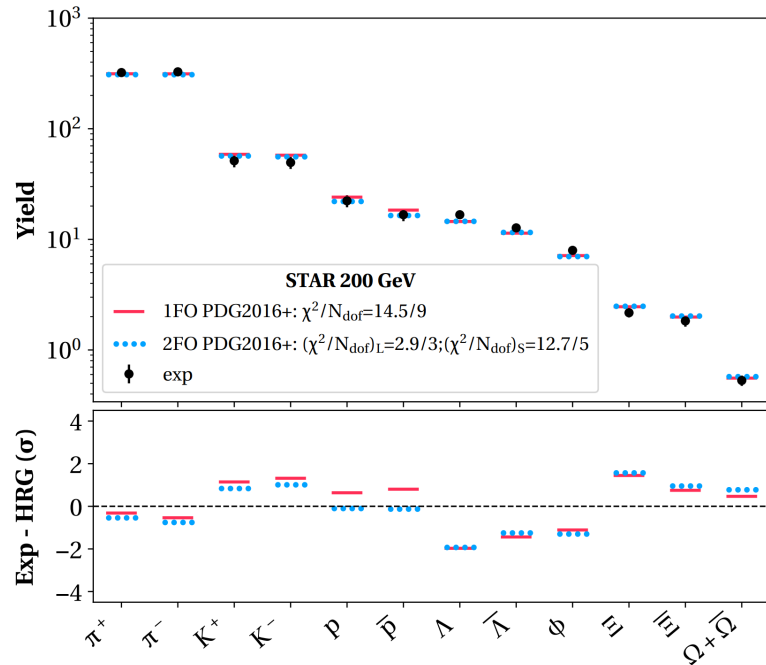


ALICE Collaboration, EPJC 84, 813 (2024)

© J. Cleymans

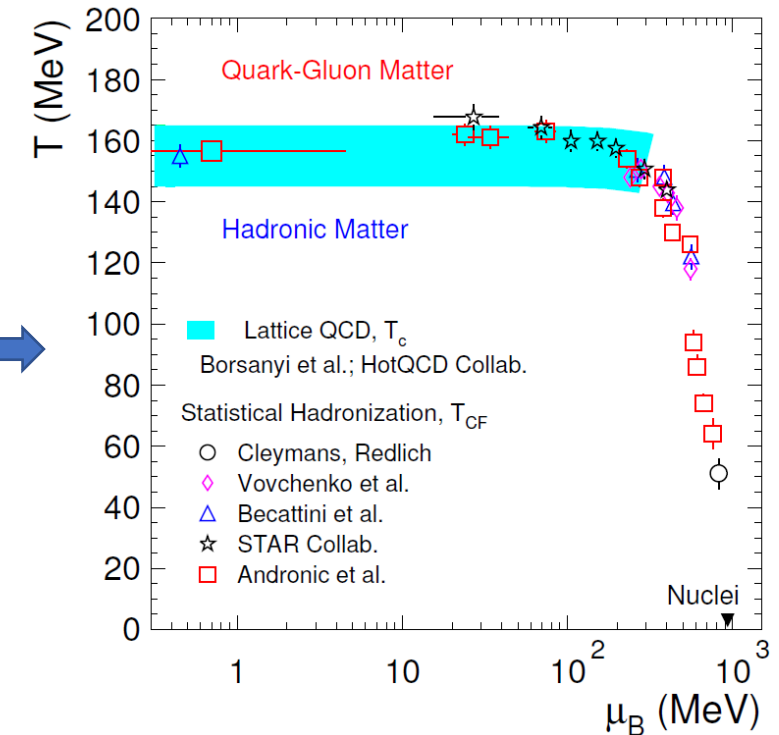
Mapping heavy-ion collisions onto the QCD phase diagram

Fit hadron yields with the HRG model



P. Alba et al. (UH group), Phys. Rev. C 101, 054905 (2020)

$$\begin{array}{c}
 T(\sqrt{s_{NN}}) \\
 \mu_B(\sqrt{s_{NN}})
 \end{array}$$

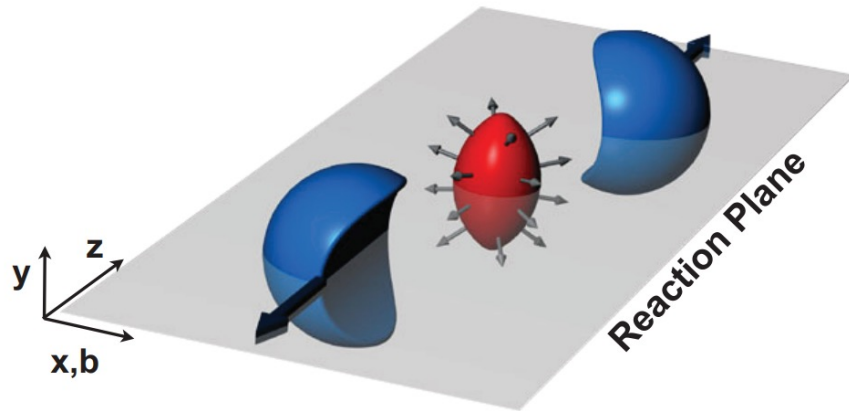


A. Andronic et al., Nature 561, 321 (2018)

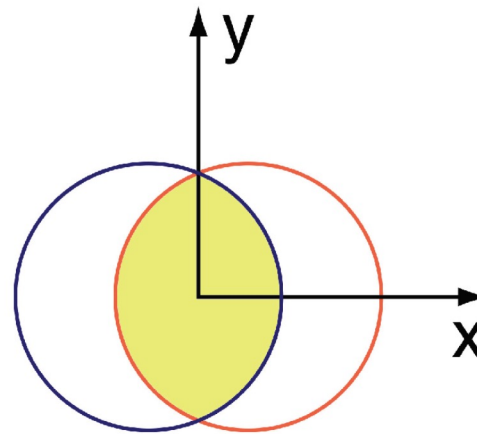
$$\sqrt{s_{NN}} \searrow \longrightarrow \mu_B \nearrow$$

Heavy-ion freeze-out probes the QCD transition at zero and non-zero baryon density

$$\frac{dN}{d\Phi dp_t} = \frac{dN}{dp_t} (1 + v_1(p_t) \cos(\Phi) + 2v_2(p_t) \cos(2\Phi) + \dots)$$



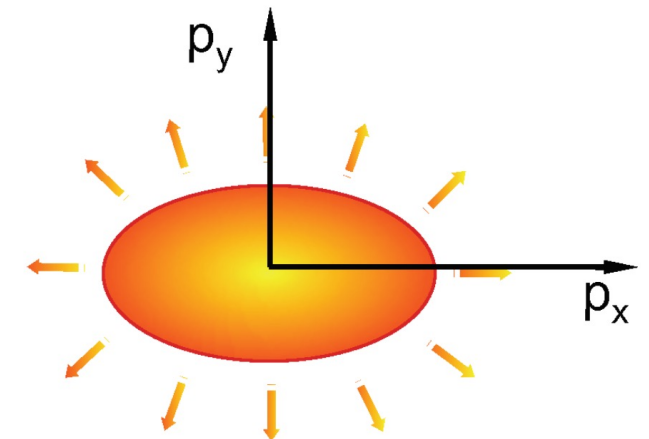
coordinate-space anisotropy



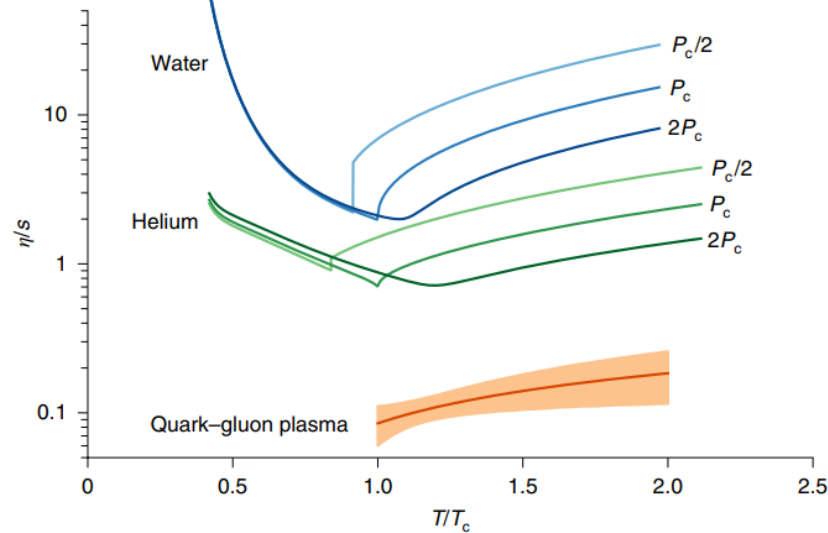
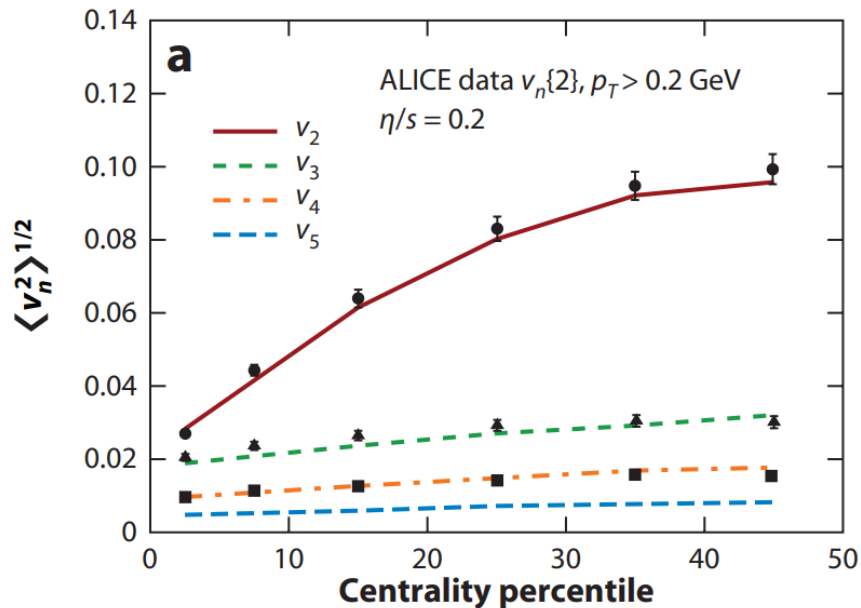
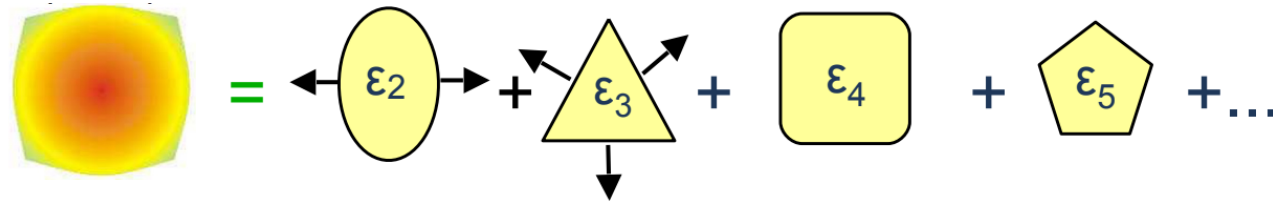
$$\partial_\mu T^{\mu\nu} = 0$$

(relativistic) hydrodynamics

momentum-space anisotropy



QGP is a fluid!



From ALICE Collaboration, Phys. Rev. Lett. 107, 032301 (2011)

From Bernhard, Moreland, Bass, Nature Physics 15, 1113 (2019)

Shear viscosity over entropy of QGP: $\eta/s \leq 0.25$ – a near ‘perfect’ fluid

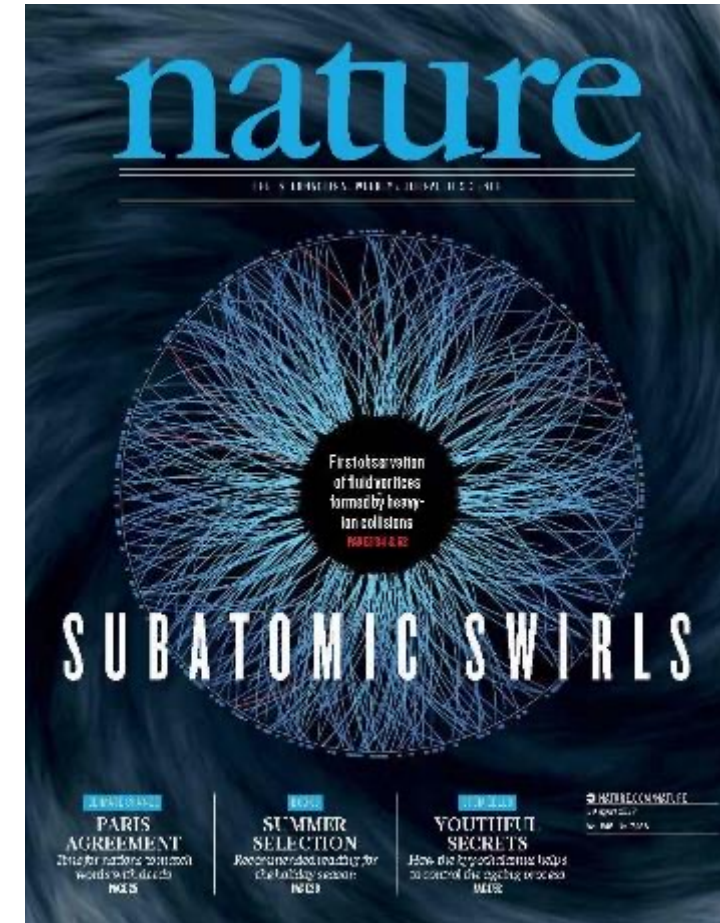
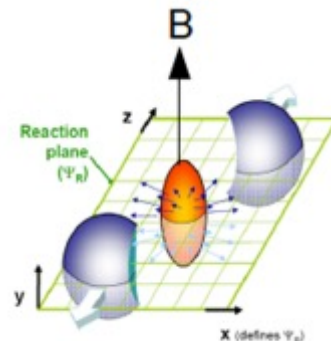
The early universe behaves like a perfect liquid rather than a gas or plasma

QGP is the hottest and most vortical fluid created on Earth

	The Earth's magnetic field	0.6 Gauss
	A common, hand-held magnet	100 Gauss
	The strongest steady magnetic fields achieved so far in the laboratory	4.5×10^5 Gauss
	The strongest man-made fields ever achieved, if only briefly	10^7 Gauss
	Typical surface, polar magnetic fields of radio pulsars	10^{11} Gauss
	Surface field of Magnetars	10^{15} Gauss

<http://solomon.as.utexas.edu/~duncan/magnetar.html>

Heavy ion collisions: the strongest magnetic field ever achieved in the laboratory
Off central Gold-Gold Collisions at 100 GeV per nucleon
 $e B(\tau=0) \sim 10^{19}$ Gauss



The challenge of discovering the QCD critical point

Disclaimer: This is my area of active research

QCD critical point

Ordinary fluid (e.g. water)

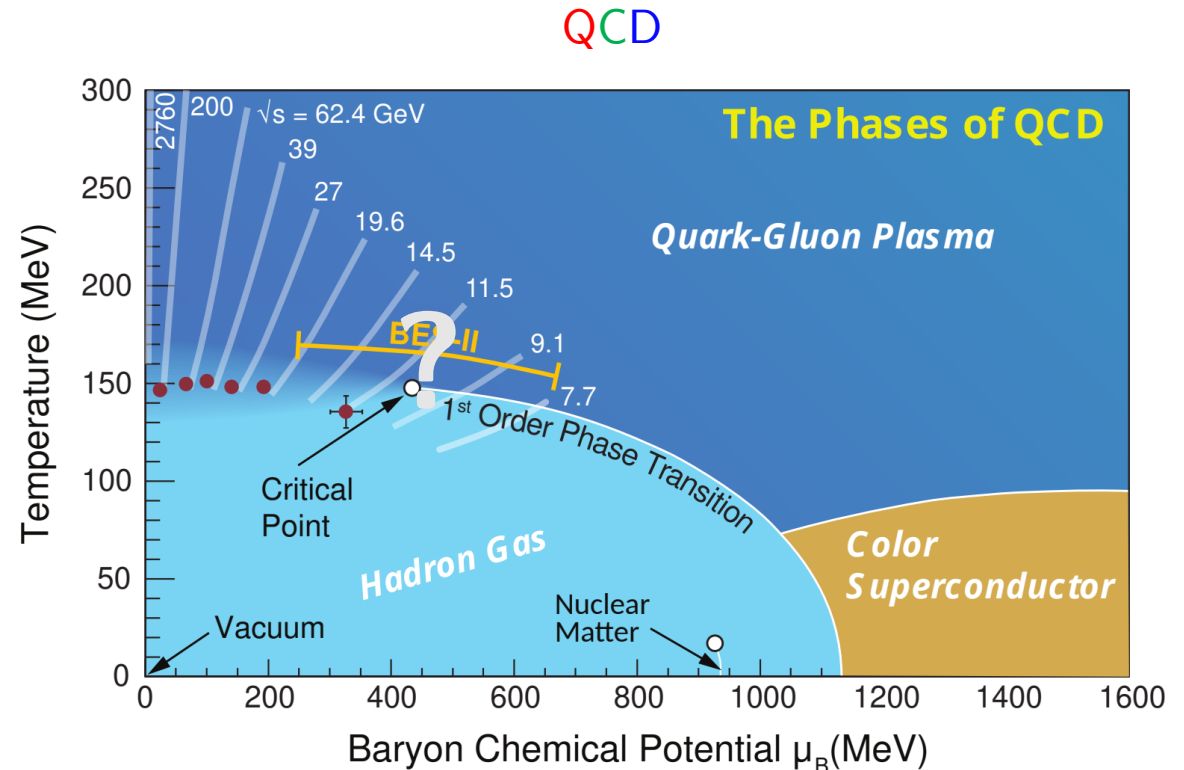
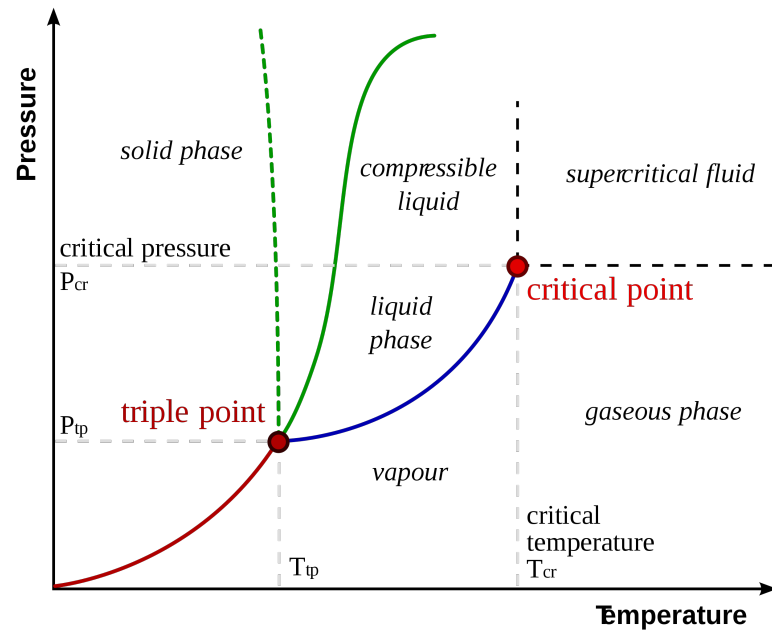


Figure from Bzdak et al., Phys. Rept. '20 and 2015 Nuclear Long Range Plan

What is the nature of the quark-hadron transition at finite baryon density?

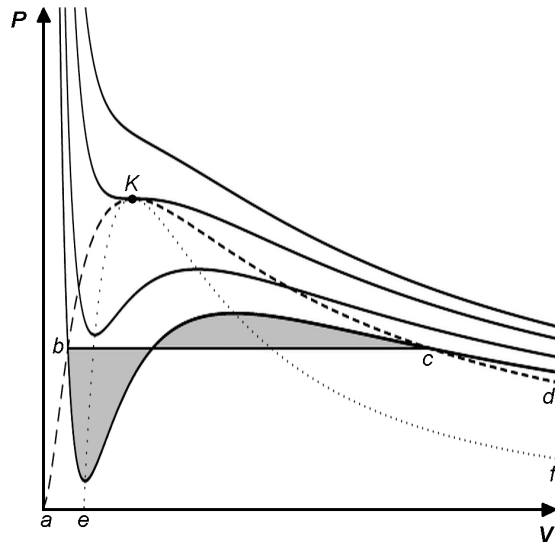
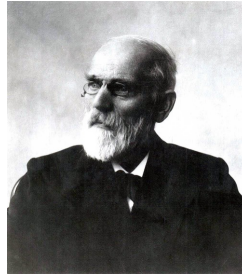
*Is there a QCD phase transition and **critical point**? Where?*

Lattice QCD: sign problem prevents simulations at non-zero baryon density

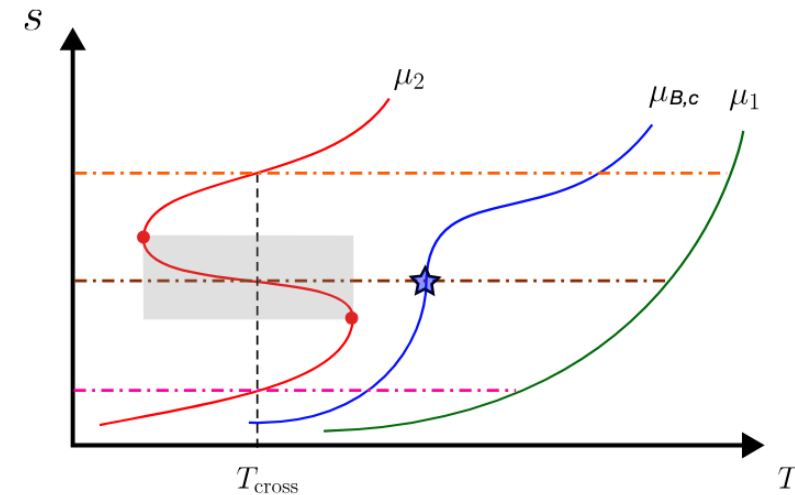
Heavy-ion collisions: access finite density but might be too short-lived to observe a signal

Extrapolating critical point from lattice

van der Waals (1873)



change the variables



Critical Point:

$$\left(\frac{\partial P}{\partial \rho_B}\right)_T = 0, \quad \left(\frac{\partial^2 P}{\partial \rho_B^2}\right)_T = 0.$$

$$\left(\frac{\partial T}{\partial s}\right)_{\mu_B} = 0, \quad \left(\frac{\partial^2 T}{\partial s^2}\right)_{\mu_B} = 0.$$

Extrapolate from $\mu_B = 0$!

Looking for entropy crossings

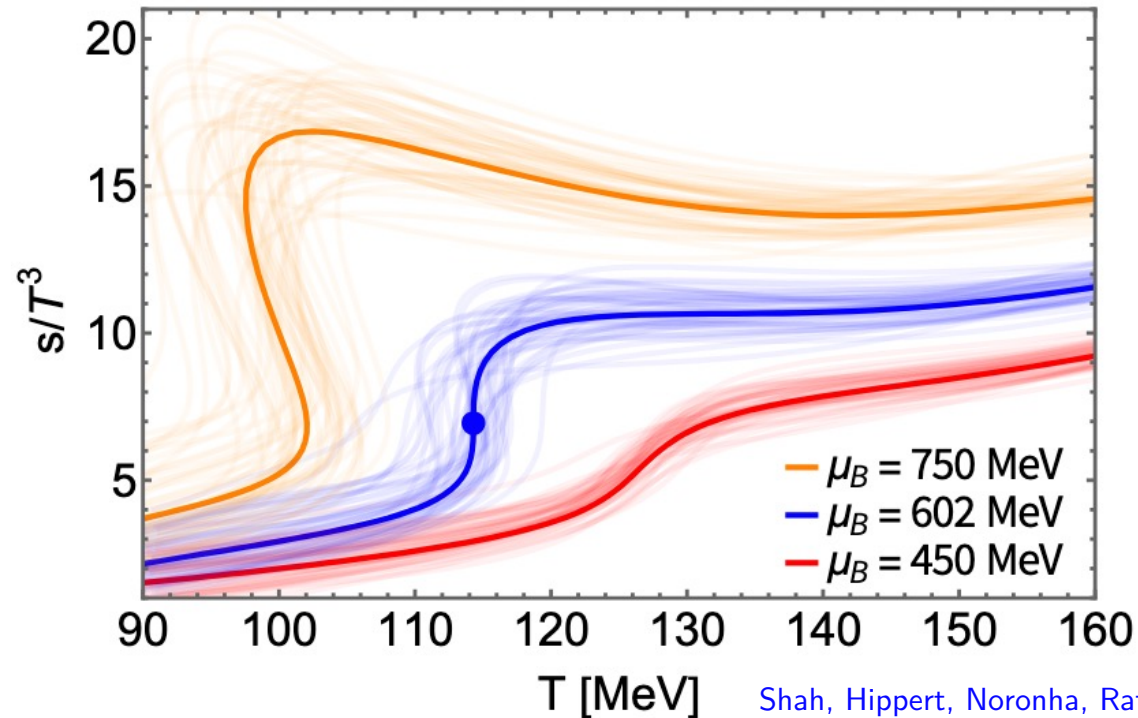
- Critical point ruled out (2σ level) at $\mu_B < 400$ MeV

Borsanyi et al., arXiv:2502.10267

- Try going further

Expansion around $\mu_B = 0$

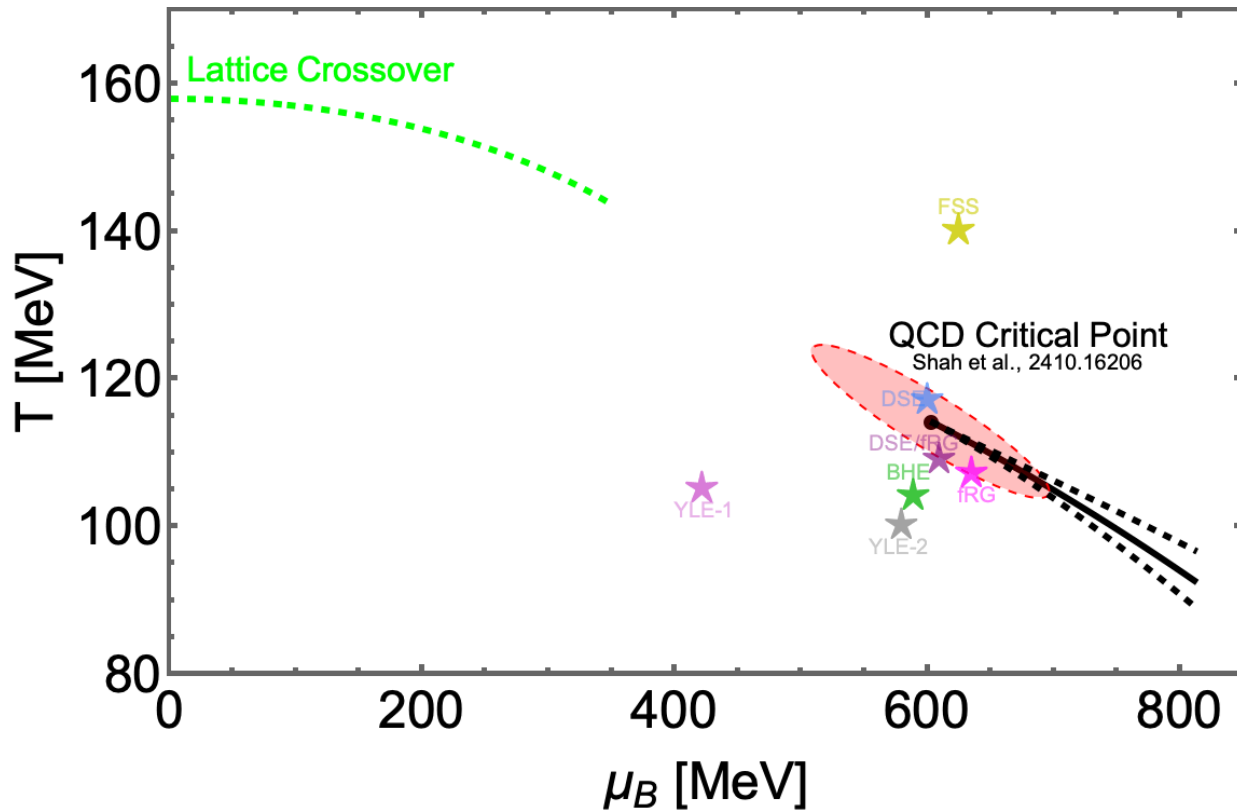
$$T_s(\mu_B; T_0) = T_0 + \alpha_2(T_0) \frac{\mu_B^2}{2}$$



Shah, Hippert, Noronha, Ratti, VV, arXiv:2410.16026

- First-order phase transition emerges at $\mu_B > 600$ MeV

QCD critical point estimates



Critical point estimate at $O(\mu_B^2)$:

$$T_c = 114 \pm 7 \text{ MeV}, \quad \mu_B = 602 \pm 62 \text{ MeV}$$

Estimates from recent literature:

YLE-1: D.A. Clarke et al. (Bielefeld-Parma), arXiv:2405.10196

YLE-2: G. Basar, PRC 110, 015203 (2024)

BHE: M. Hippert et al., arXiv:2309.00579

fRG: W-J. Fu et al., PRD 101, 054032 (2020)

DSE/fRG: Gao, Pawłowski., PLB 820, 136584 (2021)

DSE: P.J. Gunkel et al., PRD 104, 052022 (2021)

FSS: A. Sorensen et al., arXiv:2405.10278

Optimist's view: Different estimates converge onto the same region because QCD CP is likely there

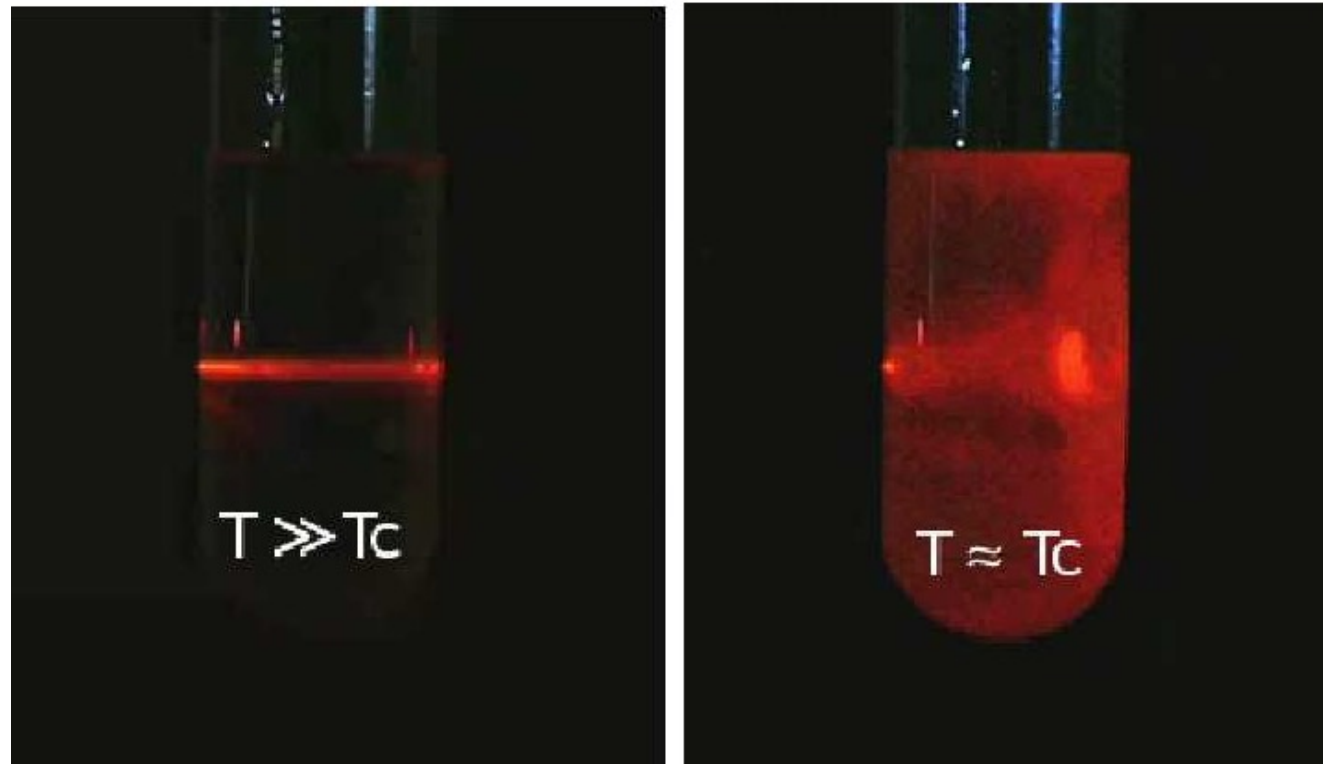
Pessimist's view: Different estimates converge onto the same region because it's the closest not yet ruled out by LQCD

Can be tested in laboratory with **heavy-ion collisions**

Critical point and fluctuations

Density fluctuations at macroscopic length scales

Critical opalescence



Unfortunately, we cannot do this in heavy-ion collisions

Event-by-event fluctuations and statistical mechanics

Consider a fluctuating number N

Cumulants: $G_N(t) = \ln \langle e^{tN} \rangle = \sum_{n=1}^{\infty} \kappa_n \frac{t^n}{n!}$

variance $\kappa_2 = \langle (\Delta N)^2 \rangle = \sigma^2$



width

skewness $\kappa_3 = \langle (\Delta N)^3 \rangle$



asymmetry

kurtosis $\kappa_4 = \langle (\Delta N)^4 \rangle - 3\langle (\Delta N^2) \rangle^2$



peak shape

Experiment:

$$P(N) \sim \frac{N_{\text{events}}(N)}{N_{\text{events}}^{\text{total}}}$$

Statistical mechanics:

Grand partition function

$$\ln Z^{\text{gce}}(T, V, \mu) = \ln \left[\sum_N e^{\mu N} Z^{\text{ce}}(T, V, N) \right],$$

$$\kappa_n \propto \frac{\partial^n (\ln Z^{\text{gce}})}{\partial (\mu_N)^n}$$

Cumulants measure chemical potential derivatives of the (QCD) equation of state

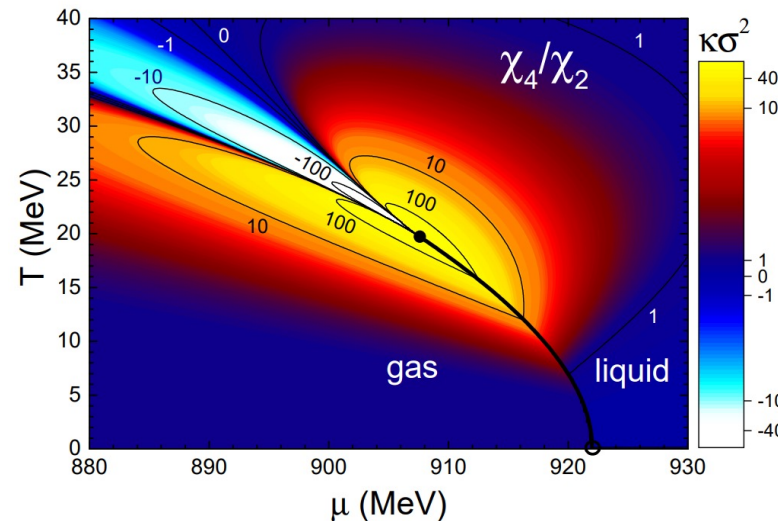
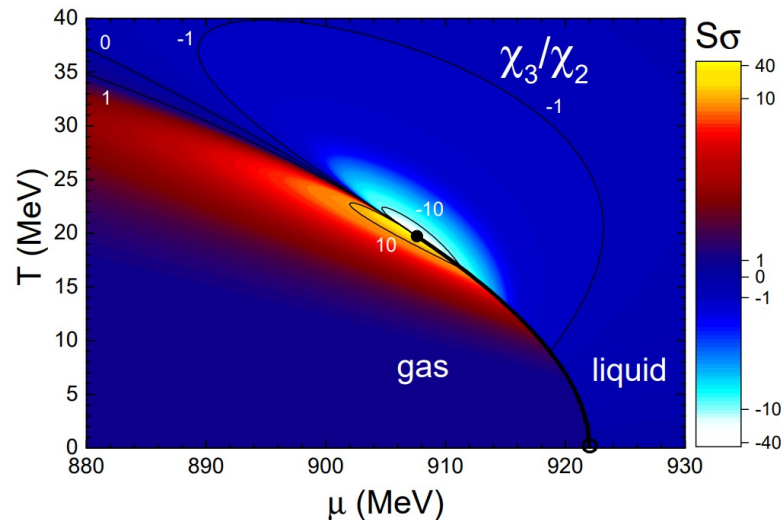
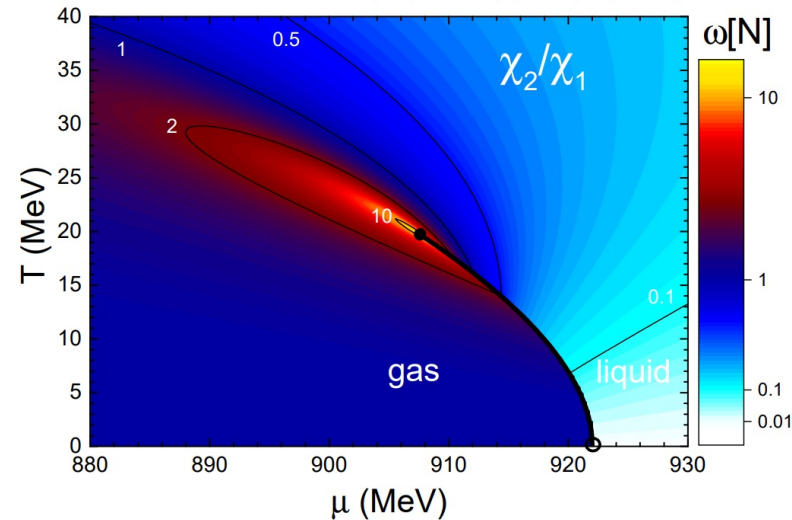
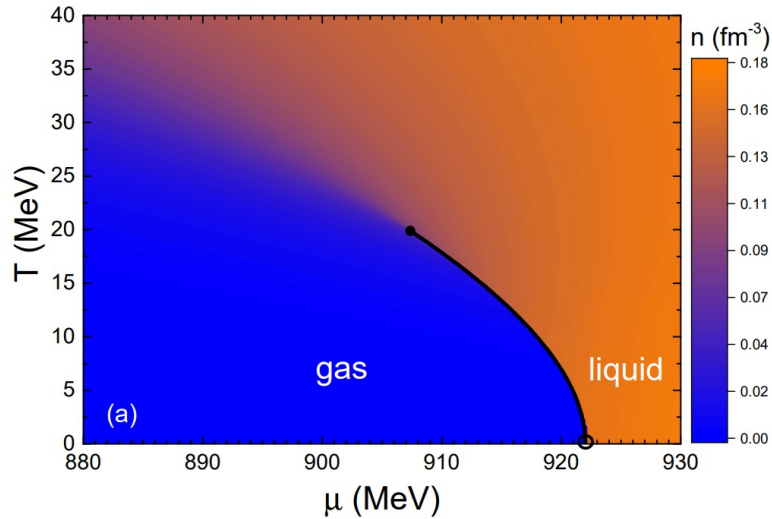
Example: (Nuclear) Liquid-gas transition

- (QCD) critical point: large correlation length and fluctuations

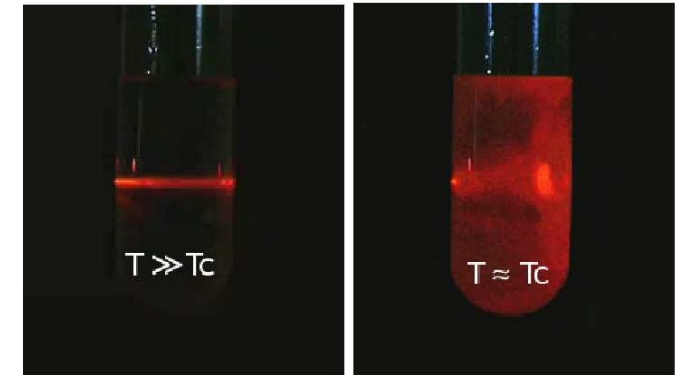
$$\kappa_2 \sim \xi^2, \quad \kappa_3 \sim \xi^{4.5}, \quad \kappa_4 \sim \xi^7$$

$$\xi \rightarrow \infty$$

M. Stephanov, PRL '09, '11



Critical opalescence



$$\langle N^2 \rangle - \langle N \rangle^2 \sim \langle N \rangle \sim 10^{23}$$

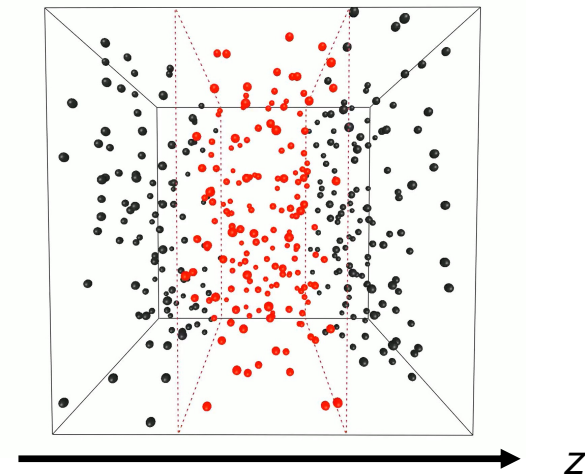
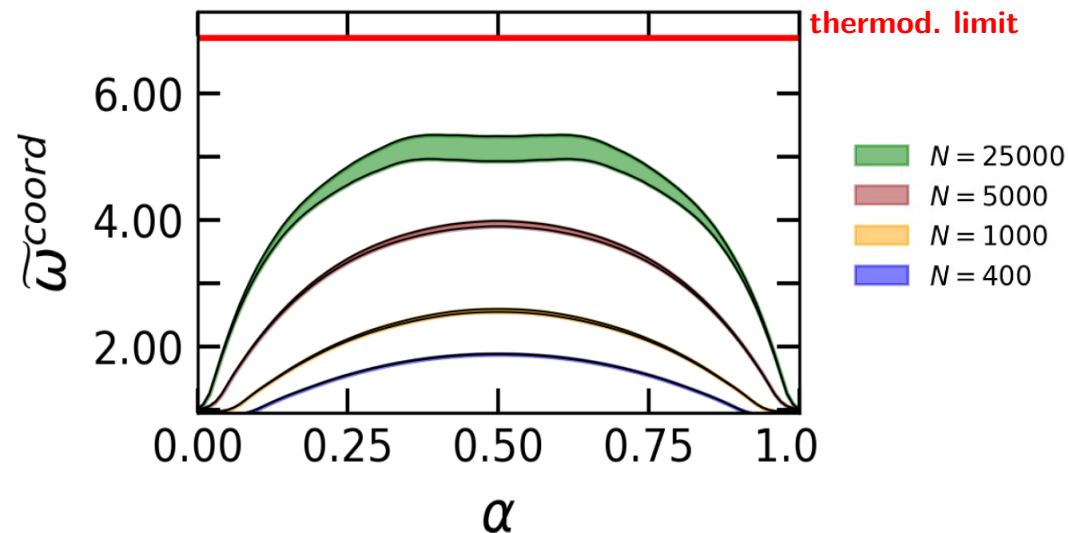
in equilibrium

Example: Critical fluctuations in microscopic simulation

V. Kuznetsov (grad student UH) et al., Phys. Rev. C 105, 044903 (2022)

Instead of observing system macroscopically, track each single particle

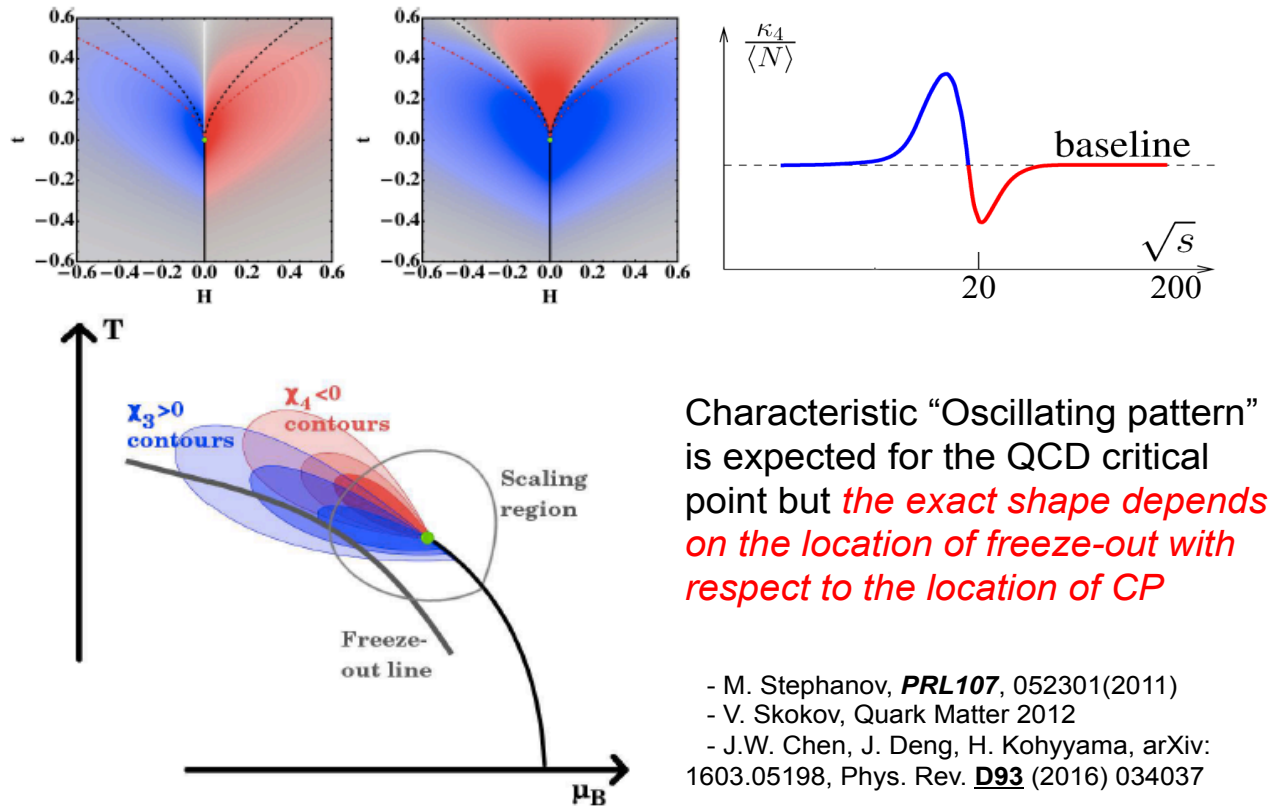
Classical molecular dynamics simulations of the **Lennard-Jones fluid** near critical point ($T \approx 1.06T_c$, $n \approx n_c$) of the liquid-gas transition



Large fluctuations survive despite strong finite-size effects and are large as advertised near the critical point

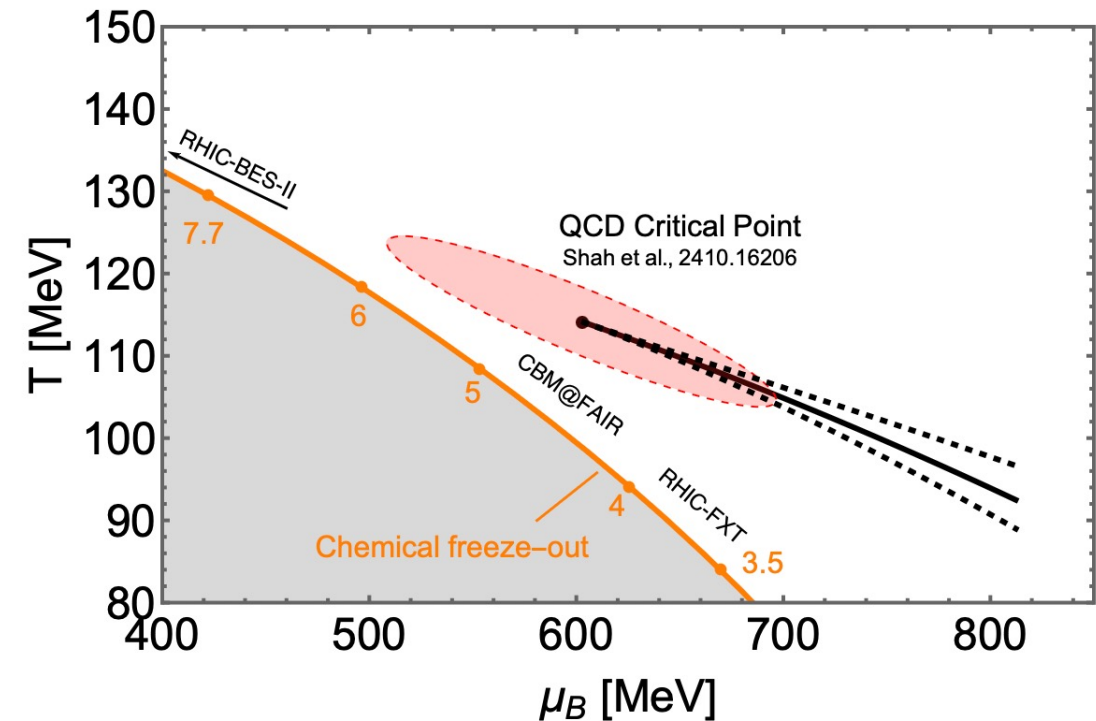
Equilibrium Expectations and Beam Energy Scan

Expectation from Calculations



N. Xu, CPOD 2016

Compare recent CP estimates and the freeze-out curve

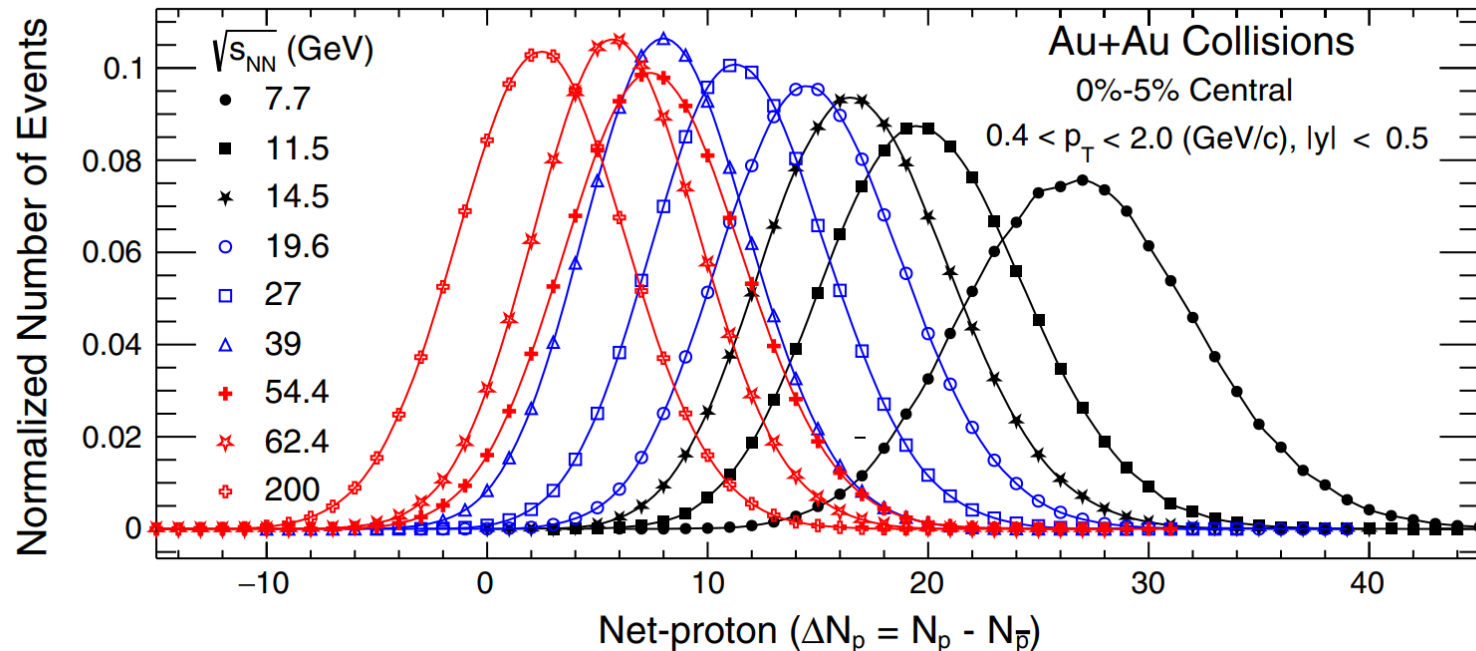


One of primary motivations for beam energy scan (BES) programs at RHIC
BES-I (7.7-200 GeV) and BES-II (3-4.5 & 7.7-39 GeV)

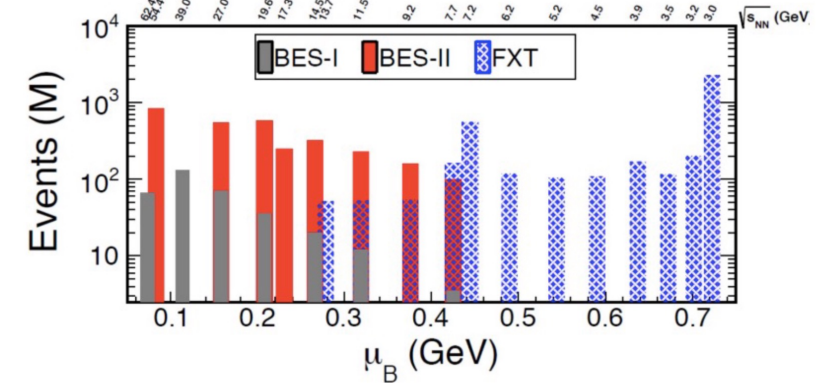
Measuring cumulants in heavy-ion collisions

Count the number of events with given number of e.g. (net) protons

STAR Collaboration, Phys. Rev. Lett. 126, 092301 (2021)



$$P(\Delta N_p) \sim \frac{N_{\text{events}}(\Delta N_p)}{N_{\text{events}}^{\text{total}}}$$



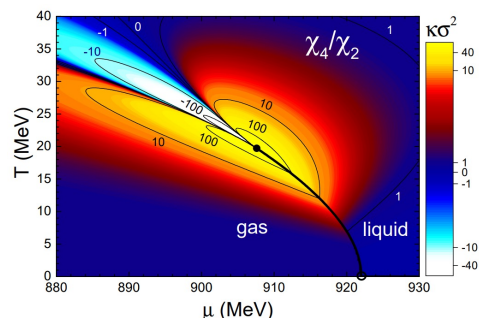
Cumulants are extensive, $\kappa_n \sim V$, use ratios to cancel out the volume

$$\frac{\kappa_2}{\langle N \rangle}, \quad \frac{\kappa_3}{\kappa_2}, \quad \frac{\kappa_4}{\kappa_2}$$

Look for subtle critical point signals

Theory vs experiment

guidance from theory

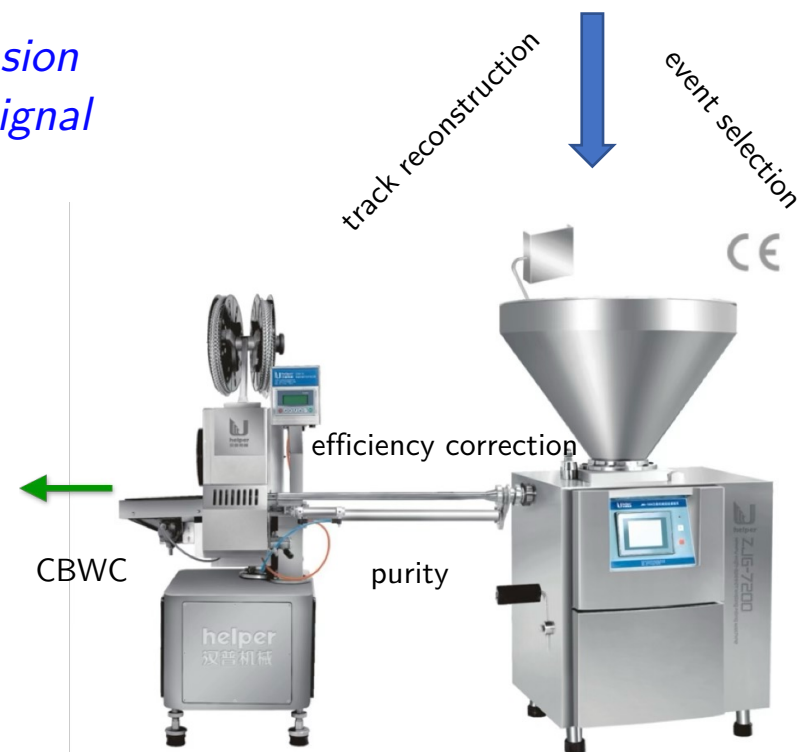
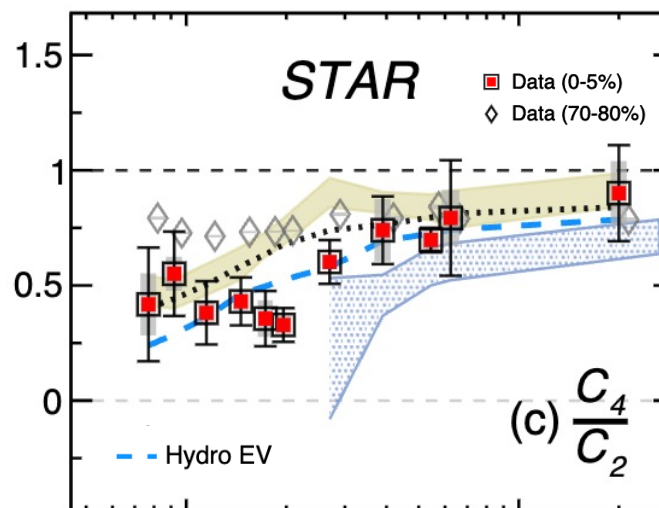
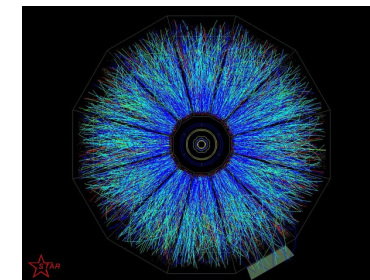


**Comparing theory and experiment
should be done very carefully**

*Some effects that are unrelated to the CP
may masquerade the signal*

*On the other hand, due to rapid expansion
of the fireball the actual critical point signal
may be very subtle*

experiment (the real thing)



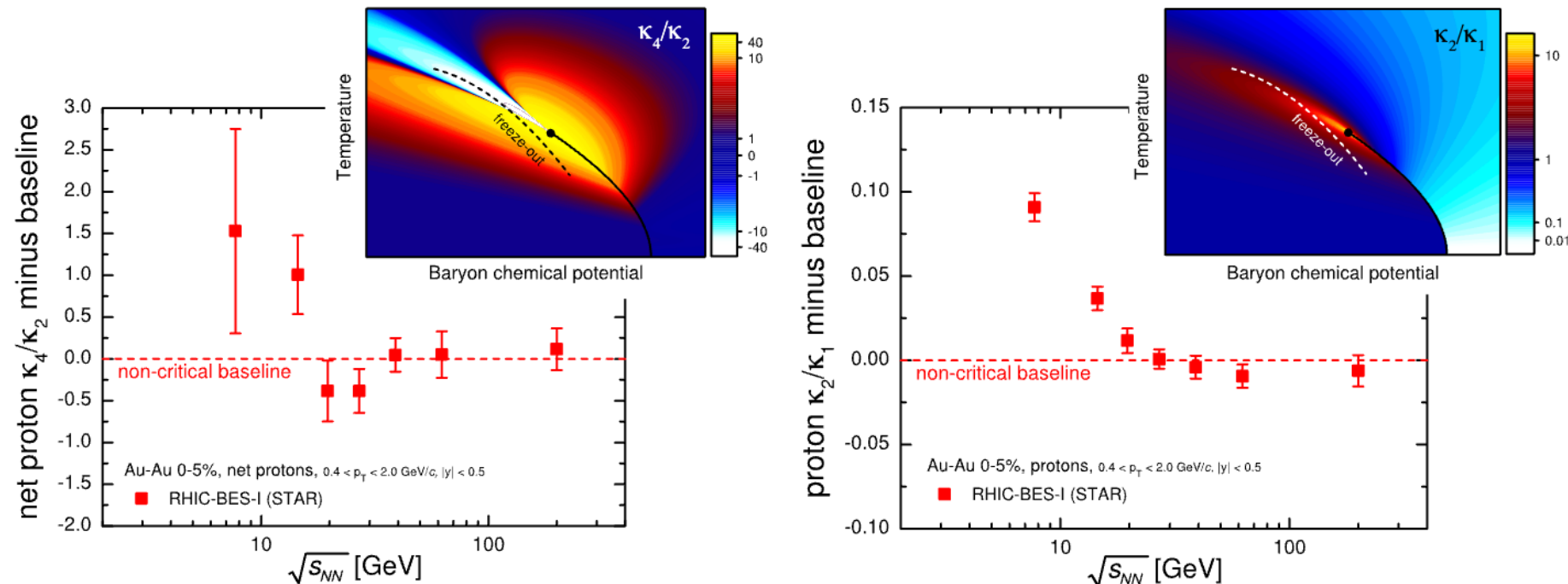
Hints from RHIC-BES-I

Quantitative calculations of critical fluctuations are still not available

State-of-the-art non-critical baseline computed using hydrodynamics

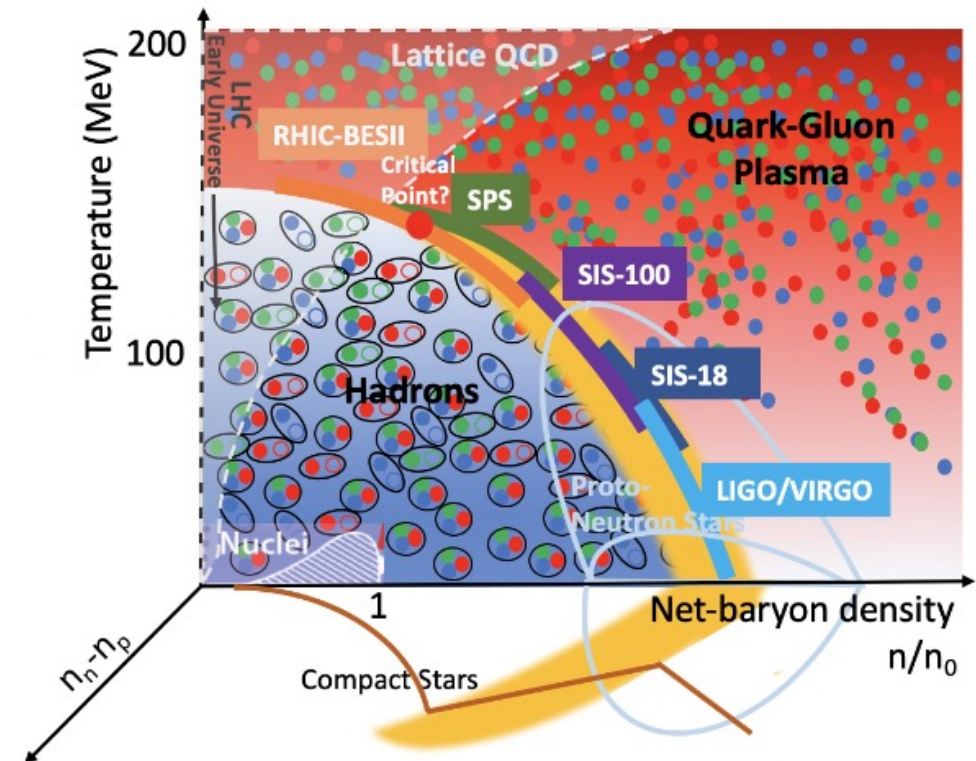
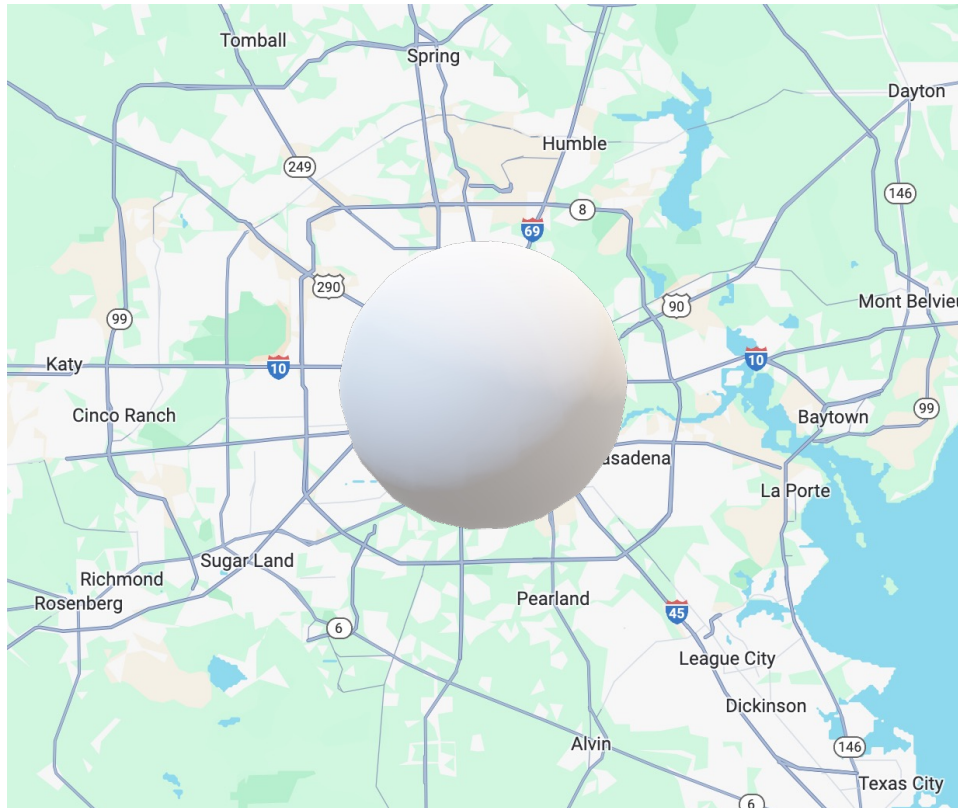
VV, V. Koch, C. Shen, Phys. Rev. C 105, 014904 (2022)

Subtract it from the data and look for a possible signal of CP



Analysis of RHIC-BES-II data in progress

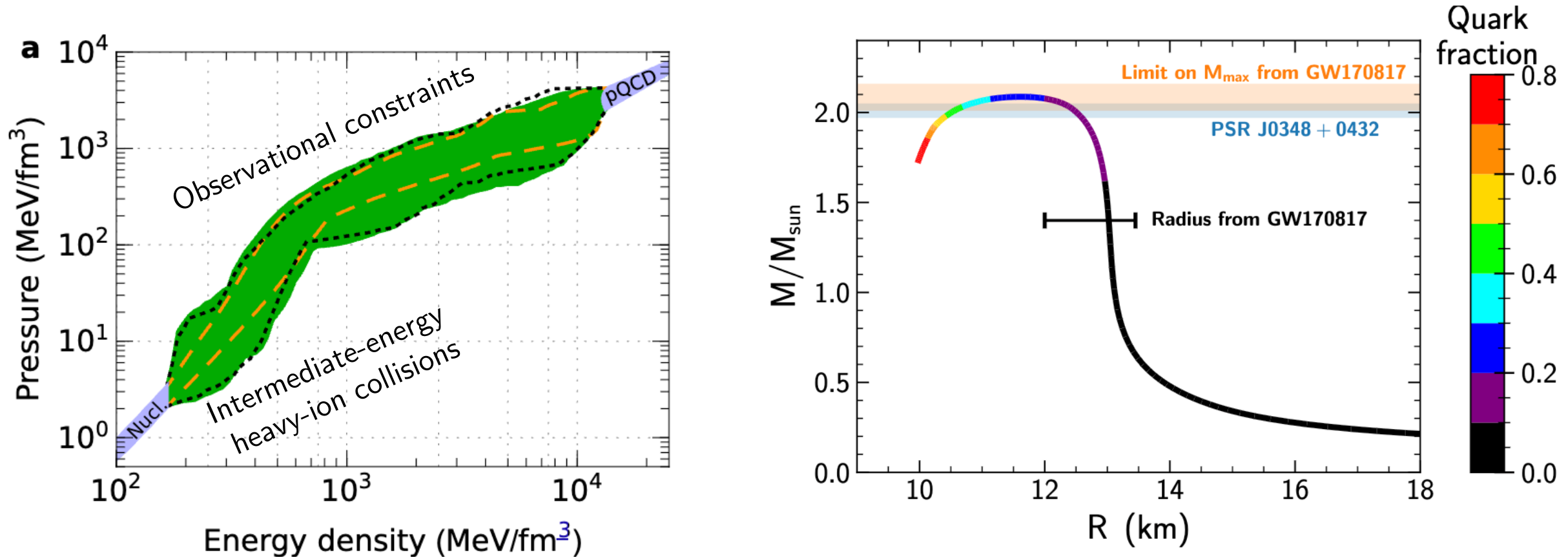
Neutron stars are extremely compact objects (1-2 solar masses confined to an 8-mile sphere)



Pressure of dense nuclear matter balances the gravitational pull

QCD in astrophysics

Properties of dense nuclear matter define how heavy neutron stars can be and how large they are



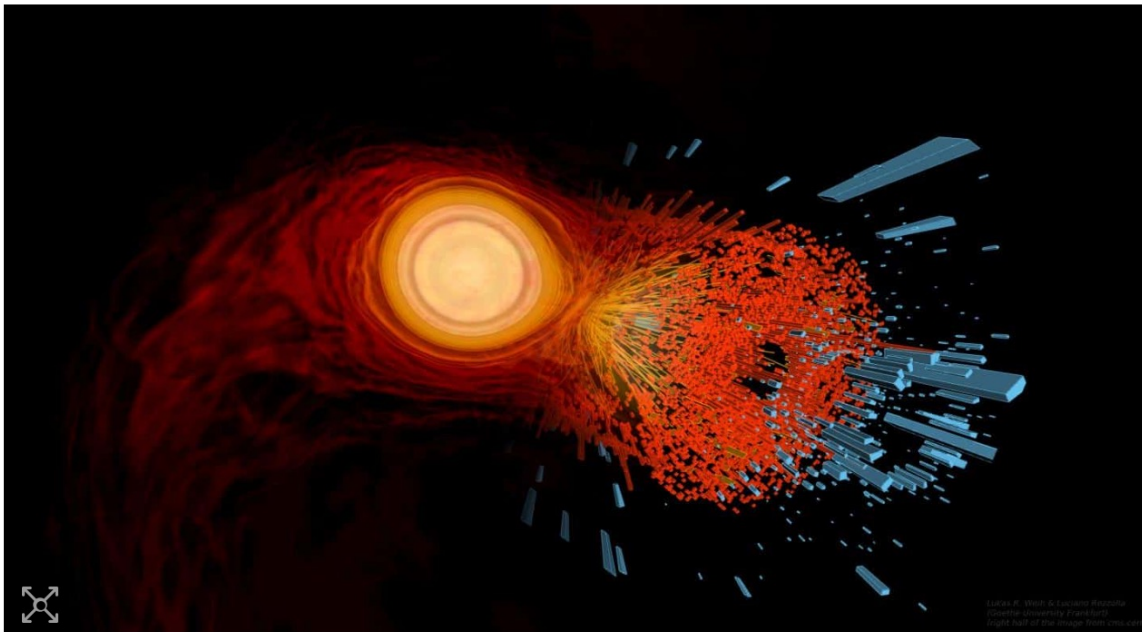
Intermediate energy heavy-ion collisions probe same dense nuclear matter

The ultimate “heavy-ion” collision

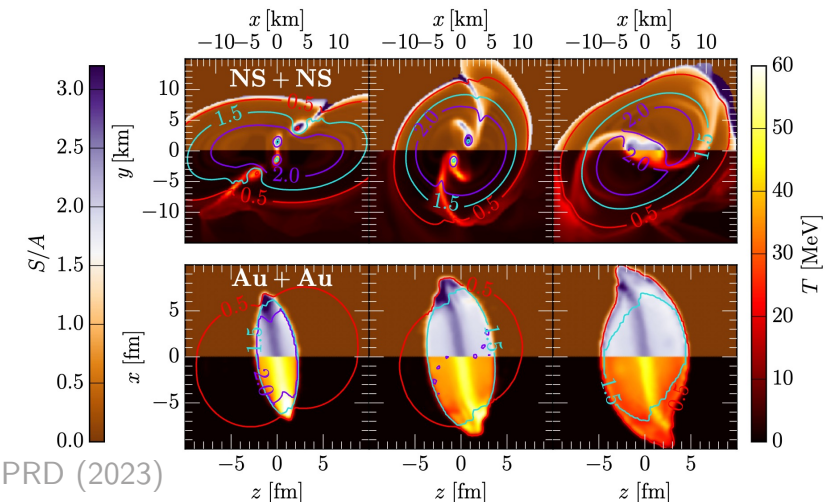
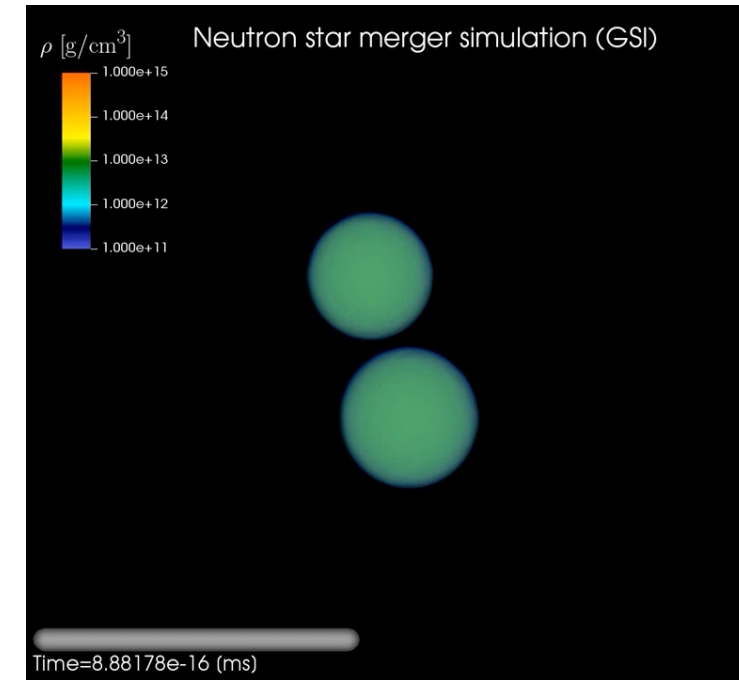
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Gravitational waves from neutron-star mergers could reveal quark-gluon plasma

15 May 2020



Dark star crashes: the computer simulation of two merging neutron stars (left) blended with an image of heavy-ion collisions at CERN to highlight the connection of astrophysics with nuclear physics. Courtesy: Lukas R Weih and Luciano Rezzolla/Goethe University Frankfurt and CMS/CERN)



Most et al., PRD (2023)